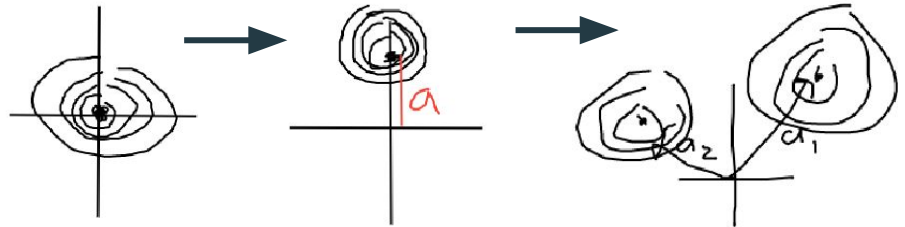


# Diffraction and Forms Factors

1. Born Approximation
2. Differential Cross Section
3. Form Factor vs Structure Function

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2\hbar^4} \left| \int d^3r \mathcal{V}(r) e^{i(\vec{k}_f - \vec{k}_i) \cdot \vec{r}} \right|^2$$

$$\begin{aligned} V(\vec{q}) &= \sum_a \int d^3r e^{i\vec{q} \cdot \vec{r}} V(\vec{r} - \vec{a}) \\ &= \sum_a e^{i\vec{q} \cdot \vec{a}} \int d^3r e^{i\vec{q} \cdot (\vec{r} - \vec{a})} V(\vec{r} - \vec{a}) \\ &= v(\vec{q}) s(\vec{q}), \end{aligned}$$



$$\frac{d\sigma}{d\Omega} = \left( \frac{d\sigma}{d\Omega} \right)_{\text{single}} \tilde{S}(\vec{q}), \quad \tilde{S}(\vec{q}) = \left| \sum_{\delta\vec{a}} e^{i\vec{q} \cdot \delta\vec{a}} \right|^2$$

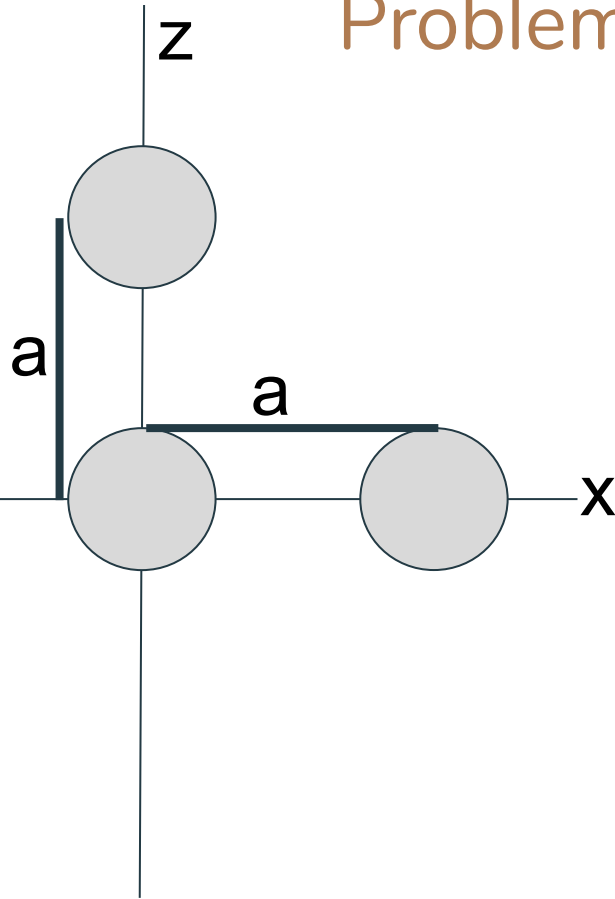
- When scatterers go from Discrete  $\rightarrow$  Continuum, you get a Form Factor

$$|s(\vec{q})|^2 = \int d\vec{a}d\vec{a}' \rho(\vec{a})\rho(\vec{a}')e^{i\vec{q}\cdot(\vec{a}-\vec{a}')}$$

- Form Factor: “form” for form of nucleus
- Structure Function: “structure” for structure of a crystal
- Independent vs Indistinguishable
- Interference Term  $\rightarrow$  Structure Function

$$\begin{aligned} |s(\vec{q})|^2 &= \sum_{a,a'} e^{i\vec{q}\cdot(\vec{a}-\vec{a}')} \\ &= N \sum_{\delta\vec{a}} e^{i\vec{q}\cdot\delta\vec{a}} \\ &= N\tilde{S}(\vec{q}), \\ \frac{d\sigma}{d\Omega} &= \frac{m^2}{4\pi^2\hbar^4} |v(q)|^2 \tilde{S}(\vec{q}). \end{aligned}$$

## Problem 7.5



The cross section of a particle with momentum  $\hbar\mathbf{k}$  off of a single target is  $\frac{d\sigma}{d\Omega} = \alpha$ ,

which is independent of  $\theta$ . Now 2 targets are placed a distance  $a$  from the first target. The incoming beam is fixed to the  $y$ - $z$  plane ( $\phi = \pi/2$ ) For what angle  $\theta$  does the differential cross section vanish?

1

$$\tilde{S}(\vec{q}) = \left| \sum_{\delta\vec{a}} e^{i\vec{q}\cdot\delta\vec{a}} \right|^2$$

$$\left( 1 + e^{-i\vec{q}\cdot\vec{a}_1} + e^{-i\vec{q}\cdot\vec{a}_2} \right)^2$$

$$\left( 1 + e^{-iq_z a} + e^{-iq_x a} \right)^2$$

$$q_z = k_z - k_z' = k(1 - \cos\theta)$$

$$q_x = \cancel{k_x} - k_x' = -k \sin\theta \cos\phi$$

$$\left| 1 + e^{\overbrace{-ik(1-\cos\theta)a}^{-iz}} + e^{ik \sin\theta \cos\phi a} \right|^2$$

$$= \left( 1 + e^z + e^{-z} + 1 + e^{-z-x} + e^z + e^x + e^{x+z} + 1 \right. \\ \left. + 2\cos z + 2\cos x - 2\cos(z+x) \right)$$

2

$$3 + 2\cos(z) + 2\cos(x) + 2\cos(z+x)$$

$$\tilde{S} \stackrel{!}{=} 0$$

Coming in from z-y plane,  $\cos\phi = 0$

$$\stackrel{\circ}{\Rightarrow} \underline{\underline{x = 0}}$$

$$z \equiv ka(1 - \cos\phi)$$

$$x \equiv ka \sin\phi \cos\phi$$

# 3

so  $\tilde{S} \Rightarrow 3 + 2\cos(z) + 2 + 2\cos(z) = 5 + 4\cos(z) = 0$

$$\Rightarrow \cos^{-1}\left\{\frac{4}{5}\right\} = z = ka(1 - \cos\phi)$$

$0 \leq \frac{2.5}{ka} \leq 2$  so there are constraints for k & a too!

$$\Rightarrow \frac{2.5}{ka} = 1 - \cos\phi \longrightarrow \phi = \cos^{-1}\left[1 - \frac{2.5}{ka}\right]$$

まっすぐまっすぐ落ちてい ← Forward, straightforward, I fall down.

# Lecture Notes

- 1 Born Approx // Foundation
- lowest order perturbation theory
  - incoming particle goes  $\vec{k} \rightarrow \vec{k}' \equiv \vec{q}$
  - single interaction w/  $V @ \vec{r}$  and its gone
  - add up all the possible point interactions + square it <sup>2</sup>
  - knowing this,  $\frac{d\sigma}{d\Omega}$  integral is INTUITIVE (Sq. Fourier transf. of  $V(\vec{r})$ )

- 2 Energy Conditions
- incoming KE  $\dot{\vec{v}}$   $\dot{V}$
  - $\vec{q} \sim$  small (unable to distinguish w/c <sup>scat. center is hit</sup>)
  - this is characteristic of elastic collisions

- 3 Origin + Translation Picture
- Origin is single scatterer Born Approx case
  - if translate by  $\vec{a}$ , potential written  $V(\vec{r}-\vec{a})$ 
    - pull out factor of  $e^{i\vec{q}\vec{a}}$  to get translated Born integral
    - separates into two terms, a phase + integ. (call  $\vec{S} + \vec{V}$ )

- 4 Expand to Multiple Scattering Centers
- there are TWO centers,  $\vec{a}_1 + \vec{a}_2$  from origin.
  - thus TWO integrals + TWO phases
    - integral dep on  $\vec{r}$ , not  $\vec{a}$ , IDENTICAL integrals
    - integral can be factored out.
  - again look @  $V + S$  terms, 'sq<sup>2</sup>' it, AND see as GEOMETRIC <sup>single</sup>  $\frac{d\sigma}{d\Omega}$  qt.

Geometric qt.  $\rightarrow$  Structure Fct.

$\frac{d\sigma}{d\Omega}$   
collection

- 5 Appreciate Simplicity of the Physics
- the hard part (integral) is only done ONCE
  - - reminiscent of Wigner Echart Theorem's efficiency
  - Structure given by Sq.<sup>2</sup> sum of exp.

- 6 Form Factors
- take  $\lim \vec{a} \rightarrow d\vec{a}$ , discrete  $\rightarrow$  continuous scatterers
  - location's become represented by a PROB. DENSITY,  $\rho(\vec{a})$ 
    - $\rho(\vec{a})d\vec{a}$ : chance that scattering center is place @  $\vec{a}$
  - we call the geometric qt. in continuous limit form factor

- 7 Lead (Pb) vs. Calcium (Ca)
- good example for eq. 5.7 // form vs struct.
  - differentiation mnemonic // Nucleus vs. Crystal
  - Form NOTE: expansion about  $\vec{q} \rightarrow 0$ 
    - take 2ND order term
    - reminiscent of a moment! (geometric form of object)
    - just like electrostatic quadrupole tensor

- 8 Independent vs. Indistinguishable
- INDEP: Scattering centers far removed from one another
    - leads to vanishing cross terms ( $e^{i\vec{q}(\vec{a}_1-\vec{a}_2)}$ ,  $e^{i\vec{q}(\vec{a}_2-\vec{a}_1)}$ ,  $e^{i\vec{q}(\vec{a}_1-\vec{a}_1)}$ )
    - b/c  $\vec{a}_1 - \vec{a}_2 \sim \text{BIG}$  and  $\vec{q} \cdot \delta\vec{a} \ll 1$
  - INDIST: momentum transf.  $\vec{q} \sim$  small
    - unable to tell w/c got hit, whole crystal moves, elastic

- 9 Interference Term (sum of all cross terms)  $\rightarrow$  Structure Fct  
AS WE WILL NOW SEE w/ EXAMPLE