

## 11. The Maxwell Equations Chapter Summary

► Three of Maxwell's equations have been explored in detail in Chapters 1 – 10,

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \rho/\epsilon_0 && \text{Gauss's Law for } \mathbf{E}(\mathbf{x}, t) \\ \nabla \cdot \mathbf{B} &= 0 && \text{Gauss's Law for } \mathbf{B}(\mathbf{x}, t) \\ \nabla \times \mathbf{E} &= -\partial\mathbf{B}/\partial t && \text{Faraday's Law}\end{aligned}$$

The fourth equation is

$$\nabla \times \mathbf{B} = \mu_0\mathbf{J} + \mu_0\epsilon_0\partial\mathbf{E}/\partial t; \quad \text{Ampère-Maxwell Law}$$

the term  $\epsilon_0\partial\mathbf{E}/\partial t$  is the displacement current density.

► For systems with polarizable materials, it is convenient to introduce the fields  $\mathbf{D}$  and  $\mathbf{H}$ . Maxwell's equations are then

$$\begin{aligned}\nabla \cdot \mathbf{D} &= \rho_{\text{free}} && \nabla \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{E} &= -\partial\mathbf{B}/\partial t && \nabla \times \mathbf{H} = \mathbf{J}_{\text{free}} + \partial\mathbf{D}/\partial t\end{aligned}$$

► We may simplify the solution of the equations formally by introducing potentials  $\mathbf{A}(\mathbf{x}, t)$  and  $V(\mathbf{x}, t)$ , in terms of which

$$\mathbf{B} = \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -\nabla V - \partial\mathbf{A}/\partial t.$$

► The field energy density and energy flux are

$$u = \frac{\epsilon_0}{2}\mathbf{E}^2 + \frac{1}{2\mu_0}\mathbf{B}^2 \quad \text{and} \quad \mathbf{S} = \frac{1}{\mu_0}\mathbf{E} \times \mathbf{B},$$

or, with linear polarizable materials,

$$u = \frac{1}{2}\mathbf{E} \cdot \mathbf{D} + \frac{1}{2}\mathbf{B} \cdot \mathbf{H} \quad \text{and} \quad \mathbf{S} = \mathbf{E} \times \mathbf{H}.$$

### ► Electromagnetic waves in vacuum

The Maxwell equations with  $\rho = 0$  and  $\mathbf{J} = 0$  have wave solutions. A harmonic plane wave has the form

$$\begin{aligned}\mathbf{E}(\mathbf{x}, t) &= \mathbf{E}_0 e^{i(\mathbf{k} \cdot \mathbf{x} - \omega t)}, \\ \mathbf{B}(\mathbf{x}, t) &= \mathbf{B}_0 e^{i(\mathbf{k} \cdot \mathbf{x} - \omega t)}.\end{aligned}$$

The field equations require that  $\mathbf{E}_0$ ,  $\mathbf{B}_0$ , and  $\mathbf{k}$  form an orthogonal triad of vectors with  $\mathbf{E}_0 \times \mathbf{B}_0$  in the direction of  $\mathbf{k}$ . Also required are

$$\omega = ck \quad \text{and} \quad B_0 = \frac{E_0}{c},$$

where  $c = 1/\sqrt{\mu_0\epsilon_0}$ . The constant  $c$  is the speed of light (and all other forms of electromagnetic radiation) in vacuum.