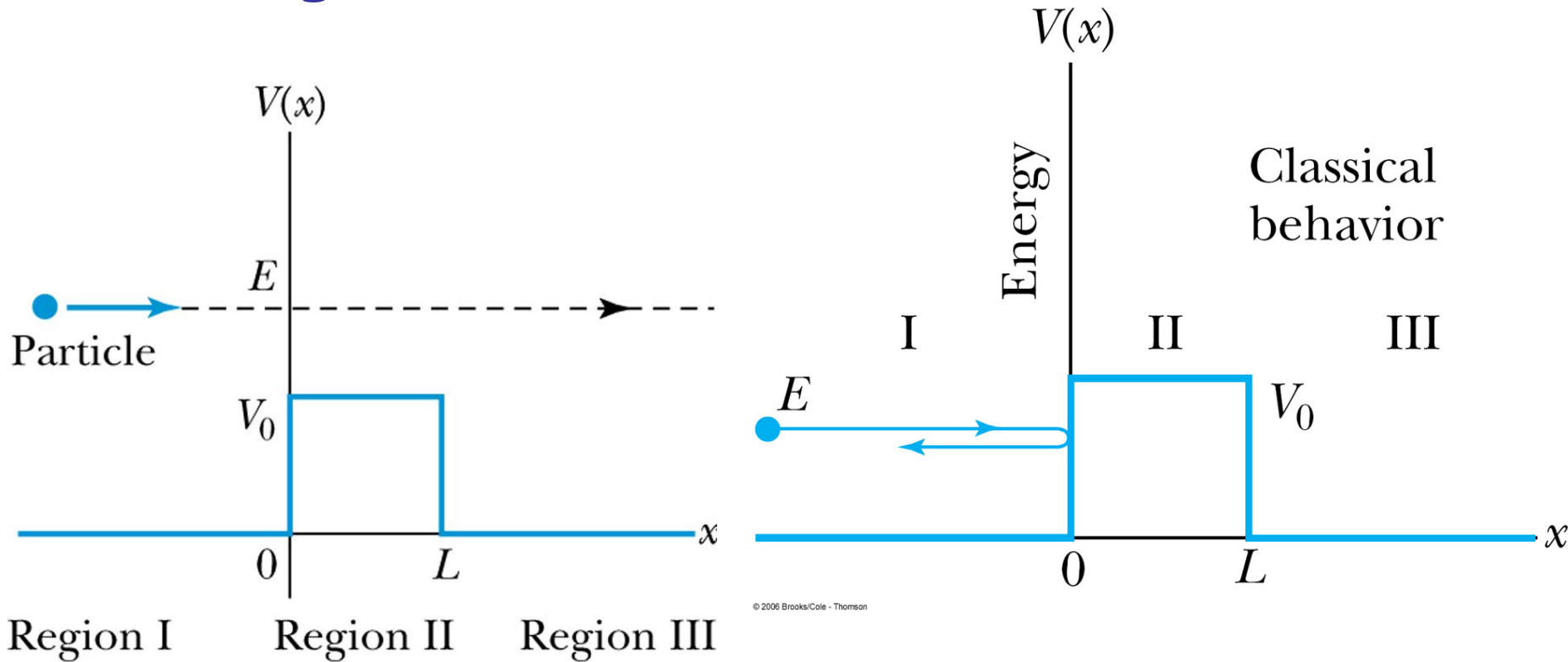
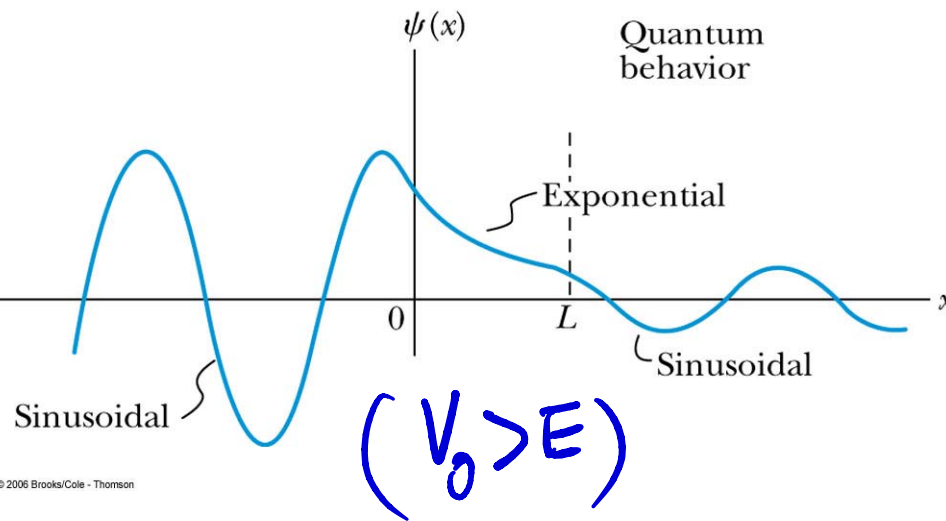


Tunneling



Classically, if a particle approaches a barrier with $E > V_0$, it is transmitted. If $E < V_0$, the classical particle will be reflected but the quantum particle can also tunnel through.

In Regions I, II and III, the Schroedinger equations are



$$e^{\pm kx}$$

$$e^{\pm ikx}$$

$$\frac{d^2 u_I}{dx^2} + \frac{2m}{\hbar^2} E u_I = 0$$

$$\frac{d^2 u_{II}}{dx^2} - \frac{2m}{\hbar^2} (V_0 - E) u_{II} = 0$$

$$\frac{d^2 u_{III}}{dx^2} + \frac{2m}{\hbar^2} E u_{III} = 0$$

$$e^{\pm ikx}$$

If we set

$$k = \sqrt{\frac{2mE}{\hbar^2}}, \quad \kappa = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}},$$

the solutions are familiar:

$$u_I(x) = Ae^{ikx} + Be^{-ikx}$$

$$u_{II}(x) = Ce^{-\kappa x} + De^{\kappa x}$$

$$u_{III}(x) = Fe^{ikx} + Ge^{-ikx}.$$

Remember that e^{ikx} moves right and e^{-ikx} moves left. In region III, we can set $G=0$ because there is only a transmitted wave there.

At this point, we match solutions at $x = 0$ and $x = L$, using, for example,

$$\begin{aligned} u_{II}(L) &= u_{III}(L) \\ \frac{du_{II}(L)}{dx} &= \frac{du_{III}(L)}{dx} \end{aligned}$$

with a similar expression connecting u_I and u_{II} at $x = 0$. The interesting quantity is the ratio $|F|^2/|A|^2$ that measures the tunneling probability. Solving for F , one finds

$$\frac{F}{A} = \frac{2e^{-ika}}{[2 \cosh(\kappa a) + i(\kappa/k - k/\kappa) \sinh(\kappa a)]}.$$

The tunneling probability is then

$$\frac{|F|^2}{|A|^2} = \left[1 + \frac{V_0^2 \sinh^2(\kappa L)}{4E(V_0 - E)} \right]^{-1}.$$

When κL is large, this becomes

$$\frac{|F|^2}{|A|^2} \rightarrow 16 \frac{E}{V_0} \left(1 - \frac{E}{V_0} \right) e^{-2\kappa L}.$$

Suppose an electron is accelerated through a 5 volt potential and strikes a 10 volt barrier of width 0.8 nm. What fraction of the electrons penetrate the barrier?

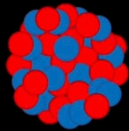
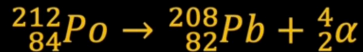
Here, $L=0.8$ nm, $V_0=10$ eV, $E=5$ eV and κ is

$$\begin{aligned}\kappa &= \frac{\sqrt{2m(V_0 - E)}}{\hbar} \\&= \frac{\sqrt{2(0.511 \times 10^6 \text{ eV}/c^2)(10 \text{ eV} - 5 \text{ eV})}}{6.58 \times 10^{-16} \text{ eV s}} \\ \kappa &= \frac{3.43 \times 10^{18} \text{ s}^{-1}}{c} = 1.15 \times 10^{10} \text{ m}^{-1}.\end{aligned}$$

From this, $\kappa L=9.2$, which is large compared to 1. We can then use

$$\frac{|F|^2}{|A|^2} = 16 \left(\frac{5 \text{ eV}}{10 \text{ eV}} \right) \left(1 - \left(\frac{5 \text{ eV}}{10 \text{ eV}} \right) \right) e^{-18.4} = 4.1 \times 10^{-8}.$$

Alpha decay of Polonium 212



$$E = 8.78\text{MeV}$$

$$T_{1/2} = 0.3\mu\text{s} = 3 \times 10^{-7}\text{s}$$

Tunnelling Probability

$$T \approx \left[\frac{16E(V_0 - E)}{V_0^2} \right] e^{-2\beta a}$$



$$T \approx e^{-2\beta a}$$

Tunnelling probability equation

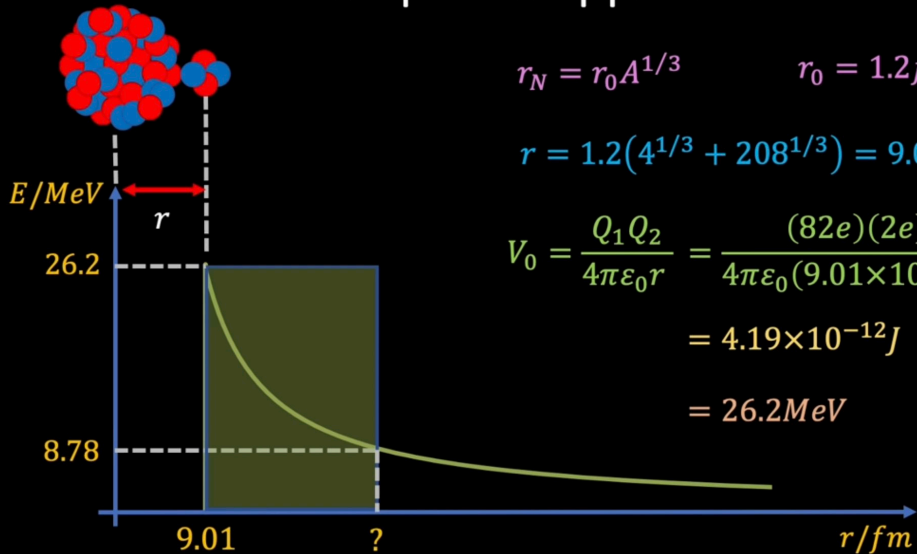


$$T = \frac{|C|^2}{|A|^2} = \frac{C^* C}{A^* A} = \left[\frac{16k^2 \beta^2}{(\beta^2 + k^2)^2} \right] e^{-2\beta a}$$

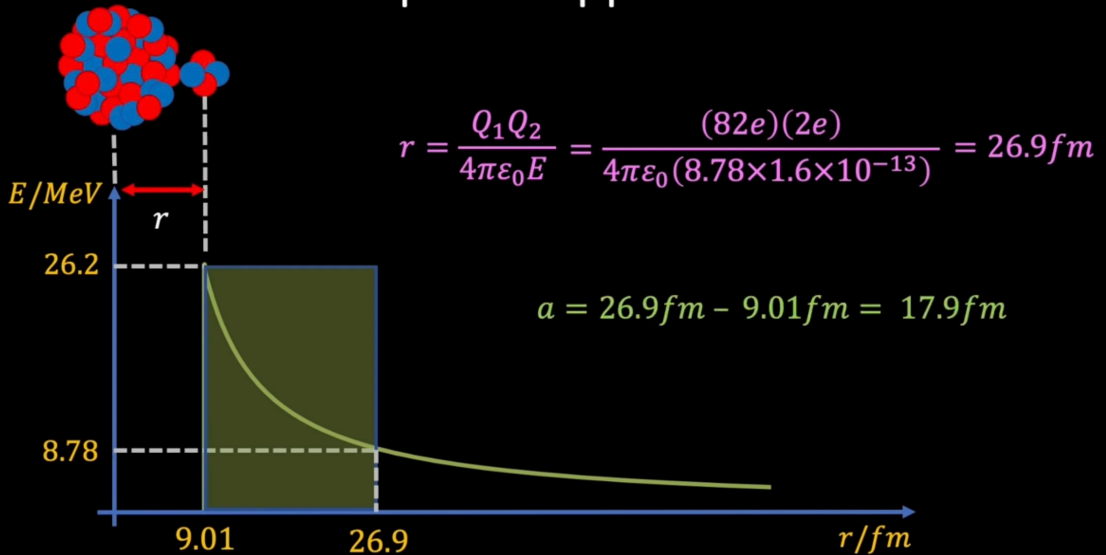
$$k = \frac{\sqrt{2mE}}{\hbar}$$

$$\beta = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$$

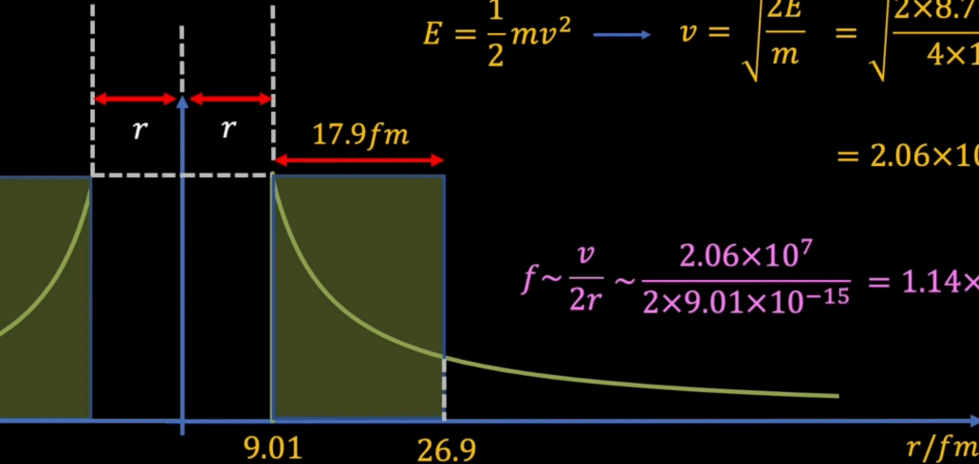
Simplified approach



Simplified approach



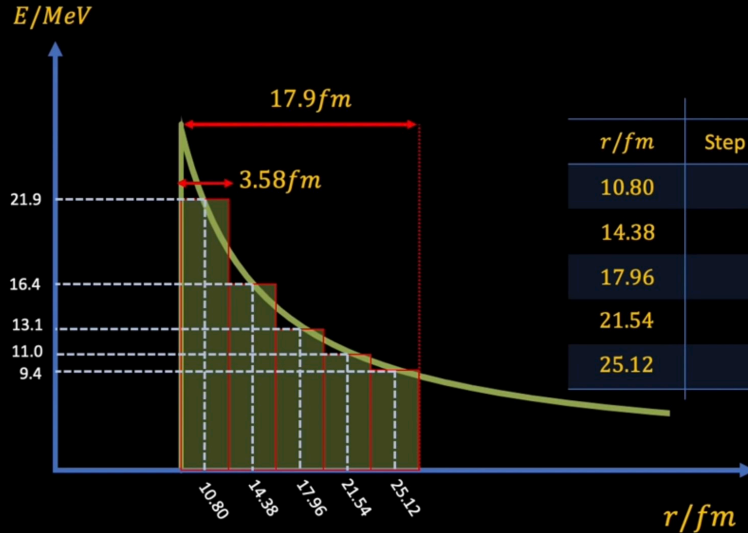
Simplified approach



$$E = \frac{1}{2}mv^2 \rightarrow v = \sqrt{\frac{2E}{m}} = \sqrt{\frac{2 \times 8.78 \times 1.6 \times 10^{-13}}{4 \times 1.67 \times 10^{-27}}} = 2.06 \times 10^7 \text{ m/s}$$

$$f \sim \frac{v}{2r} \sim \frac{2.06 \times 10^7}{2 \times 9.01 \times 10^{-15}} = 1.14 \times 10^{21} \text{ s}^{-1}$$

A better model

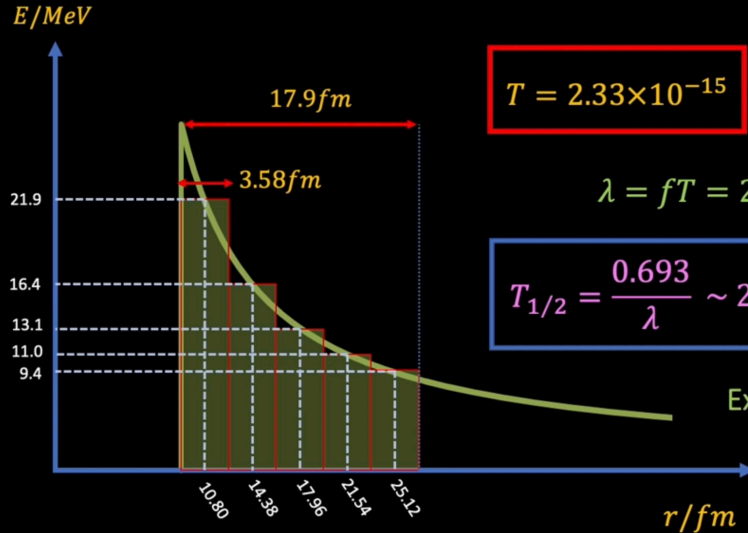


$$T = e^{-2\beta a}$$

r/fm	Step Height/MeV	Tunnelling Probability
10.80	21.9	1.16×10^{-5}
14.38	16.4	1.73×10^{-4}
17.96	13.1	1.47×10^{-3}
21.54	11.0	9.33×10^{-3}
25.12	9.4	8.46×10^{-2}

$$T = 2.33 \times 10^{-15}$$

A better model



$$T = 2.33 \times 10^{-15}$$

$$f = 1.14 \times 10^{21} \text{s}^{-1}$$

$$\lambda = fT = 2.66 \times 10^6 \text{s}^{-1}$$

$$T_{1/2} = \frac{0.693}{\lambda} \sim 2.6 \times 10^{-7} = 0.26 \mu\text{s}$$

Experiment = 0.3 μs

Not bad!

A better model

