

Atomic Physics

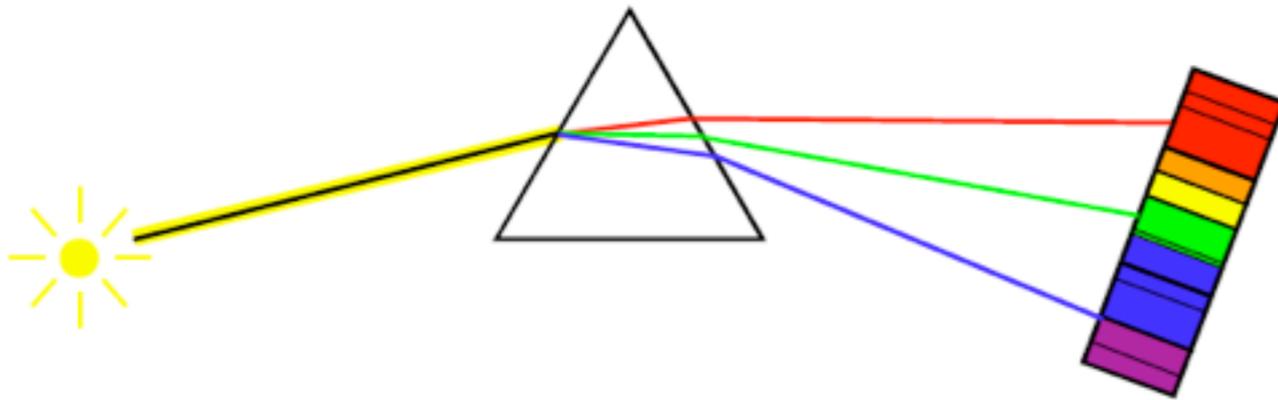
PHY 215
Thermodynamics and
Modern Physics

Spring 2026
MSU

Outline

- Spectral Lines
 - Lyman, Balmer, Paschen Series
- Models of Atoms
 - Rutherford Scattering
- Classical Hydrogen Atom
- Bohr Atom
 - Limitations

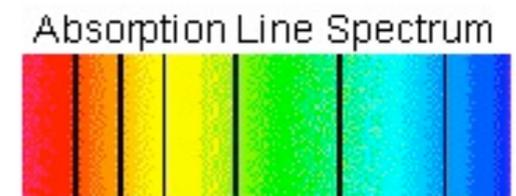
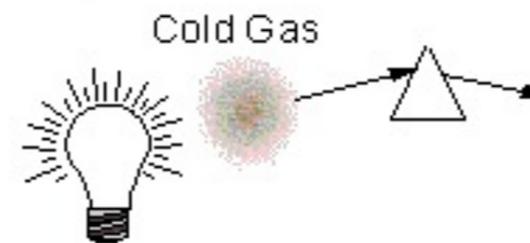
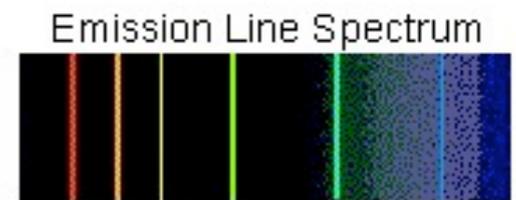
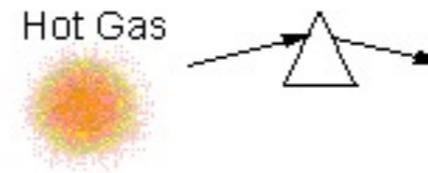
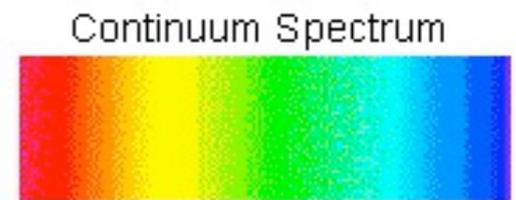
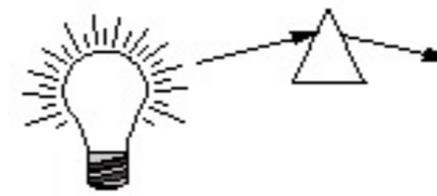
Spectral Lines

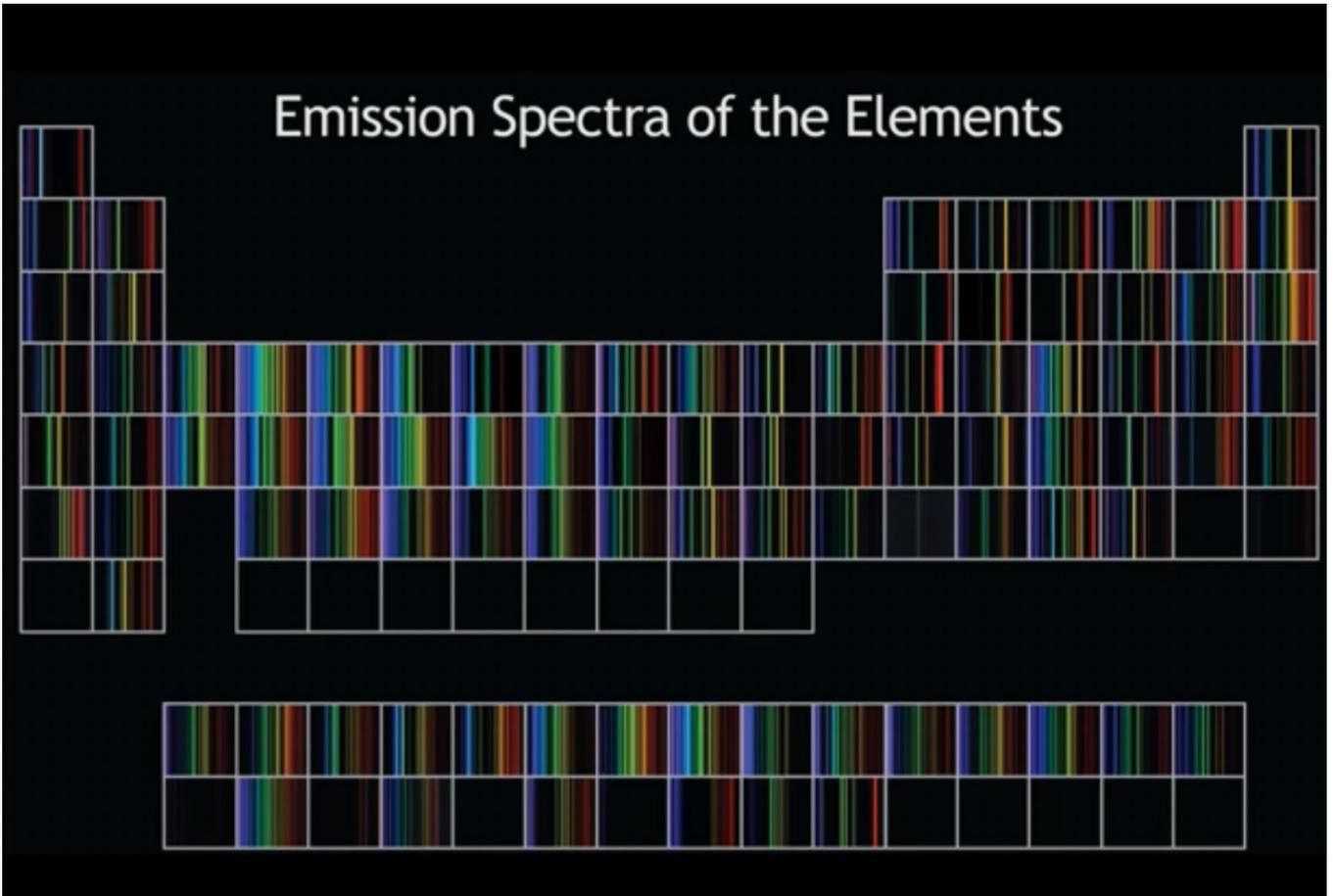
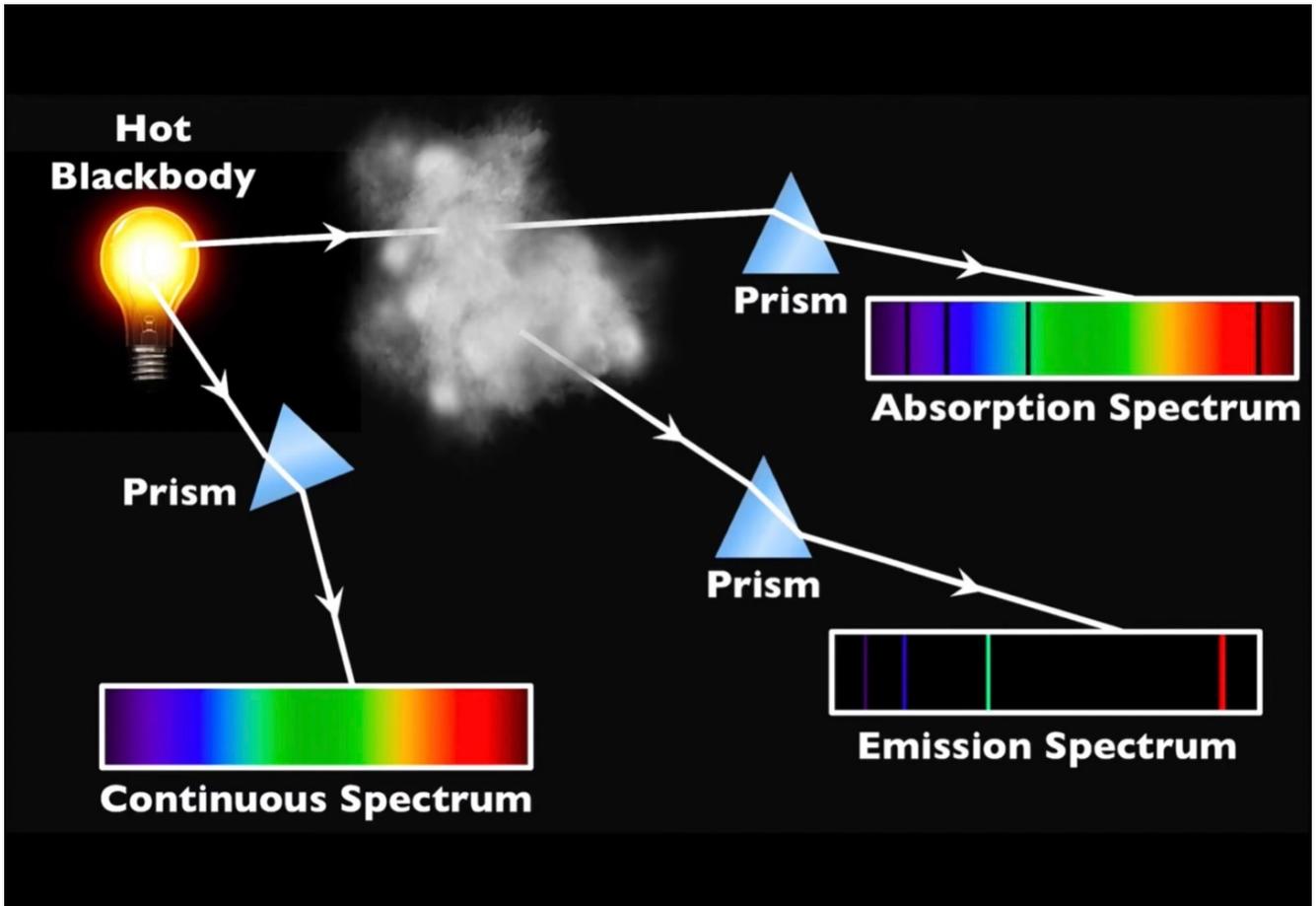


- 1814-1824: Von Fraunhofer discovered absorption lines in sun.

- 1850's: Kirchhoff discovered characteristic emission lines of elements

- 1859: Kirchhoff and Bunsen discovered new elements, Cesium and Rubidium, by first observing their spectral lines.





Concept Test

- A bright star shines through a dark gaseous nebula. The spectrum of this star will consist primarily of
 - Bright Lines
 - Dark Lines ← Star = “light bulb”
 - Neither

Balmer Series

- 1885: Balmer found a formula for the wavelengths of the spectral lines in Hydrogen:

$$\frac{1}{\lambda} = R_H \left(\frac{1}{4} - \frac{1}{n^2} \right)$$

where $n=3,4,5,\dots$

The constant R_H is Rydberg's constant:

$$R_H = 1.096776 \times 10^7 \text{ m}^{-1}$$

Rydberg Equation

Table 3.2 Hydrogen Series of Spectral Lines

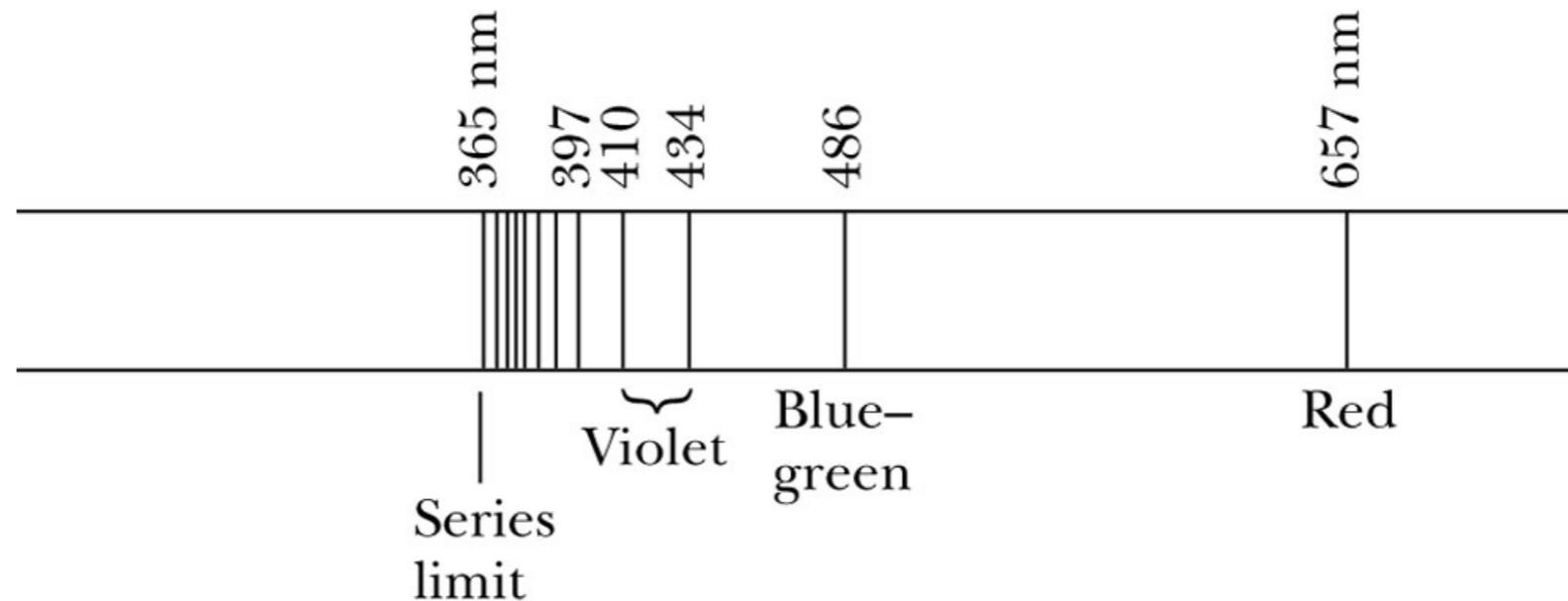
Discoverer (year)	Wavelength	n	k
Lyman (1916)	Ultraviolet	1	>1
Balmer (1885)	Visible, ultraviolet	2	>2
Paschen (1908)	Infrared	3	>3
Brackett (1922)	Infrared	4	>4
Pfund (1924)	Infrared	5	>5

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$$\frac{1}{\lambda} = R_H \left(\frac{1}{n^2} - \frac{1}{k^2} \right) \quad R_H = 1.096776 \times 10^7 \text{ m}^{-1}$$

Image: Thornton and Rex

Interpretation (Bohr)

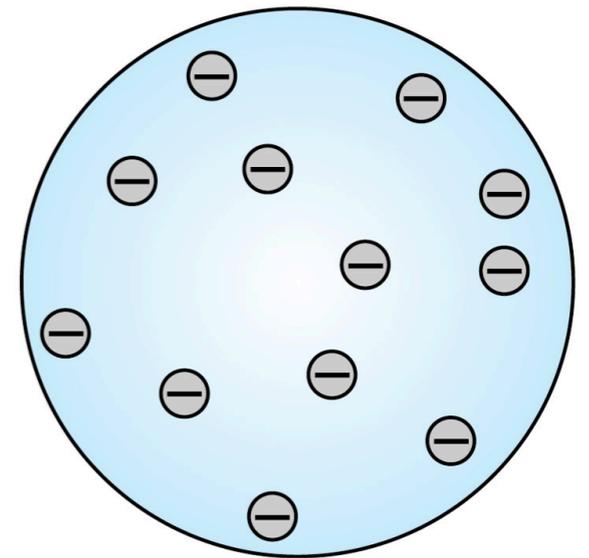


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- Discrete spectral lines
 - $E=h\nu$ (Planck)
 - **Discrete Energies!**
- Atom interacting with light
- Hence, discrete atomic “states”
- **Quantization**

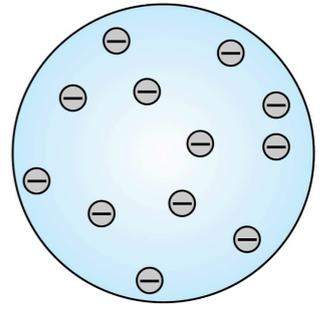
Models of Atoms

- Electrons + some positive charge must reside inside atom: but how?
- Pre-history: Thomson's Plum-Pudding
 - Electrons embedded in uniform + background (“pudding”)
- Heat atoms, electrons vibrate, create E&M radiation (light)



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Concept Test



- α particles are the nuclei of Helium atoms, have a charge of +2 and a mass of approximately 8000 times m_e . If the α particles scatter off of a “plum pudding” atom, a continuous distribution of positive charge with a few (light) electrons embedded in it, we expect:

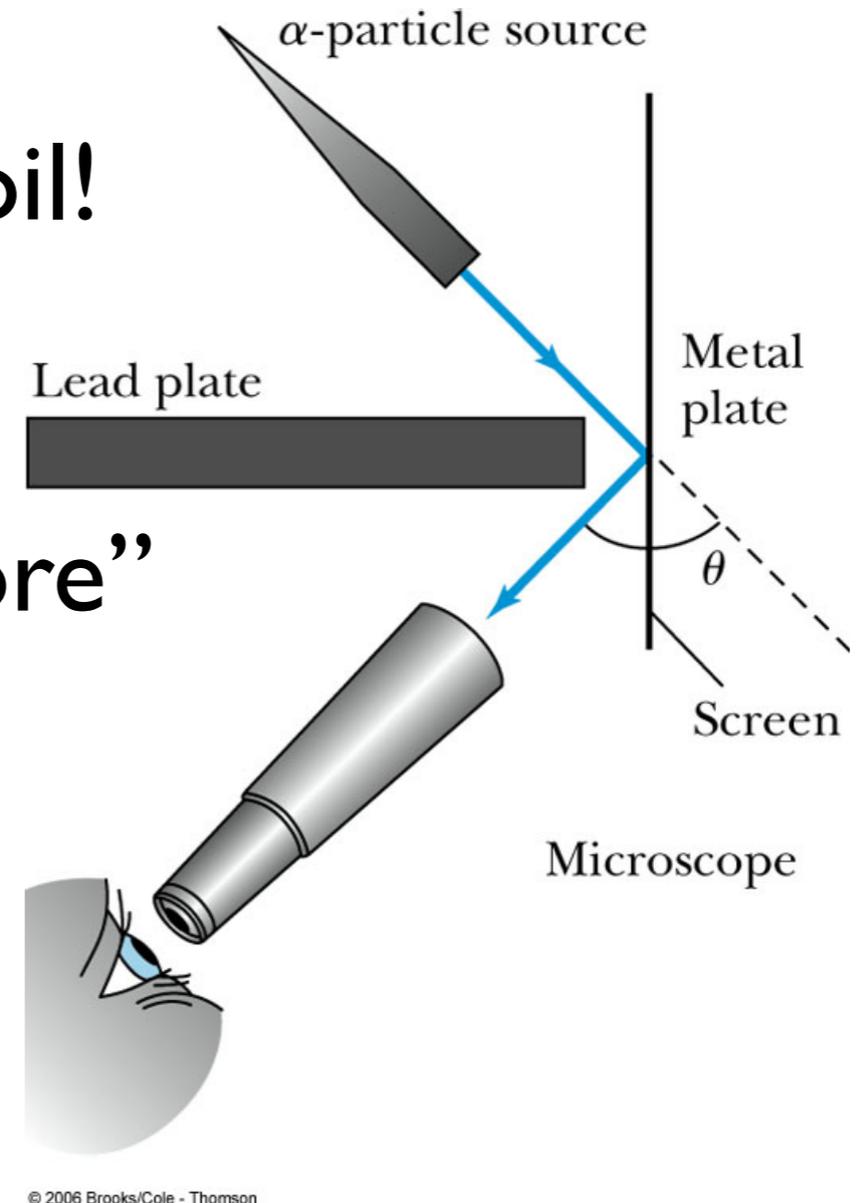
- All of the α particles to be absorbed
- All of the α particles to bounce backward
- None of the α particles to bounce backward ←

Geiger and Marsden (1909)

- α rays (modern: He nuclei)
“back-scatter” off of a thin gold foil!
- Rutherford (1911)
 - Atoms have a “hard charged core”
of size $\sim 10^{-14}\text{m}$!



Ernest Rutherford
1871-1937
Nobel Prize 1908

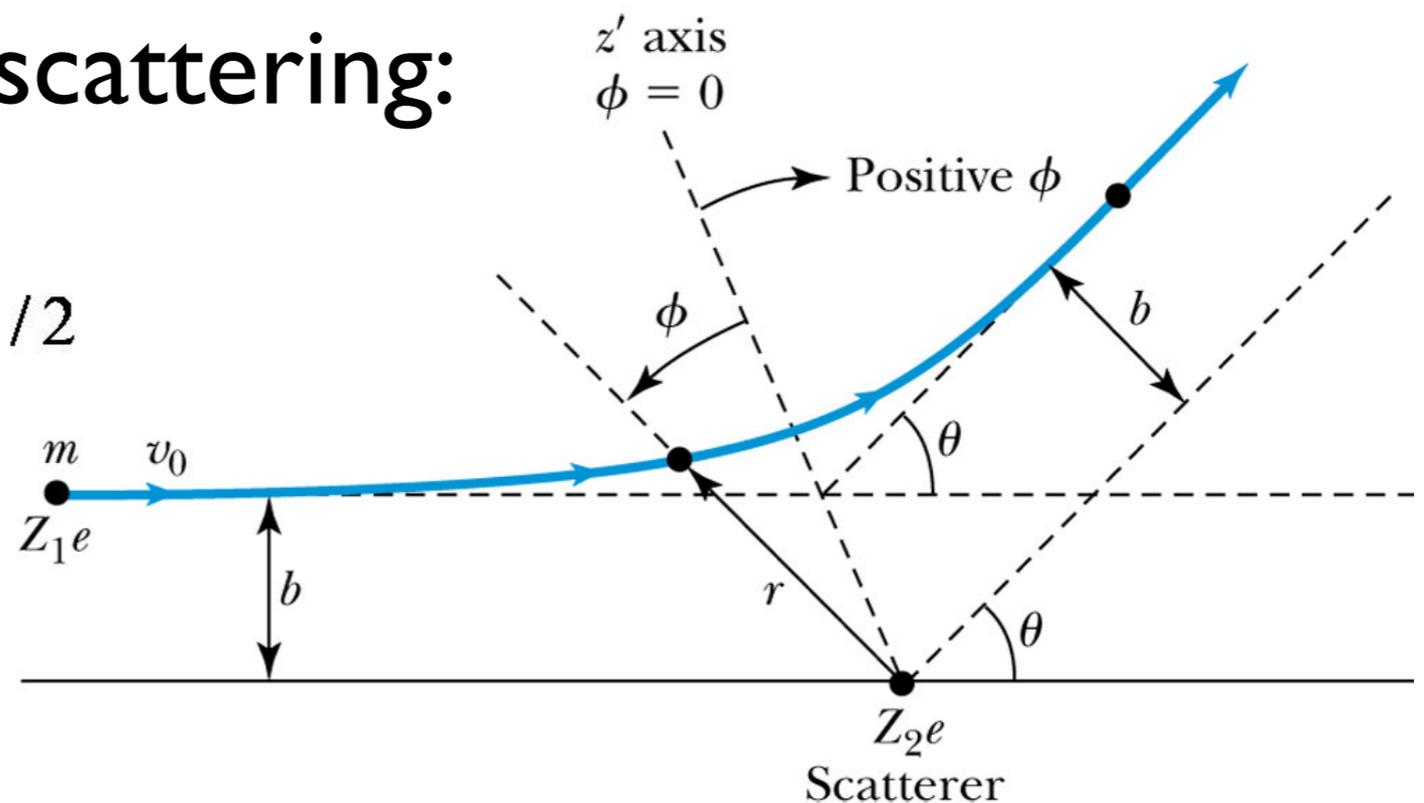


Images: Thornton and Rex
<http://wikipedia.org>

Scattering Experiments

- We study the properties of atoms (and smaller objects) using *scattering experiments*.
- Two parameters: impact parameter (b) and scattering angle (θ)
- For Coulomb scattering:

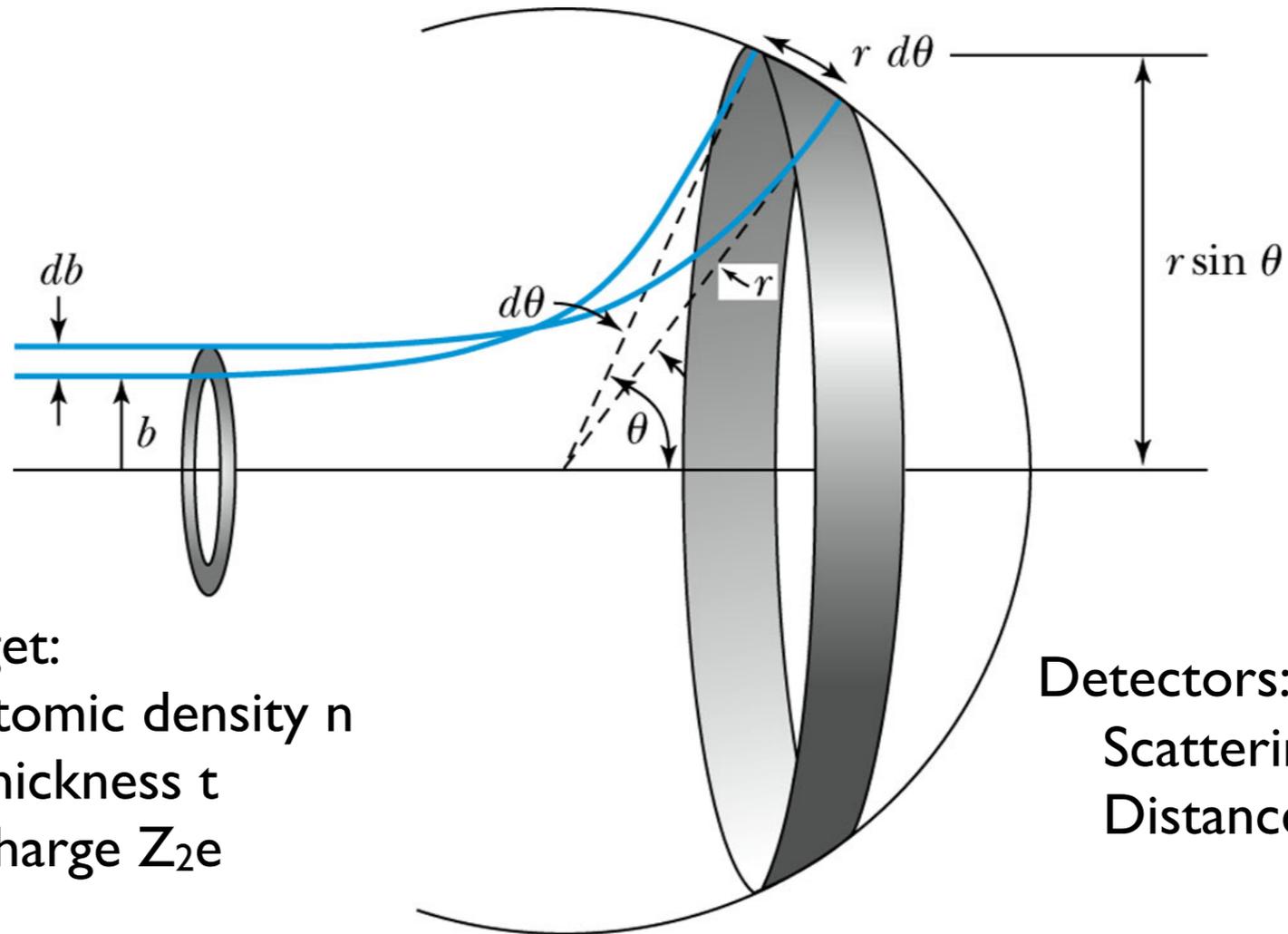
$$b = \frac{Z_1 Z_2 e^2}{8\pi\epsilon_0 K} \cot \frac{\theta}{2} \quad \text{where } K = mv_0^2 / 2$$



Images: [Thornton and Rex](#)

Rutherford Scattering Formula

N_i incoming particles
 Charge $Z_1 e$
 Kinetic Energy K



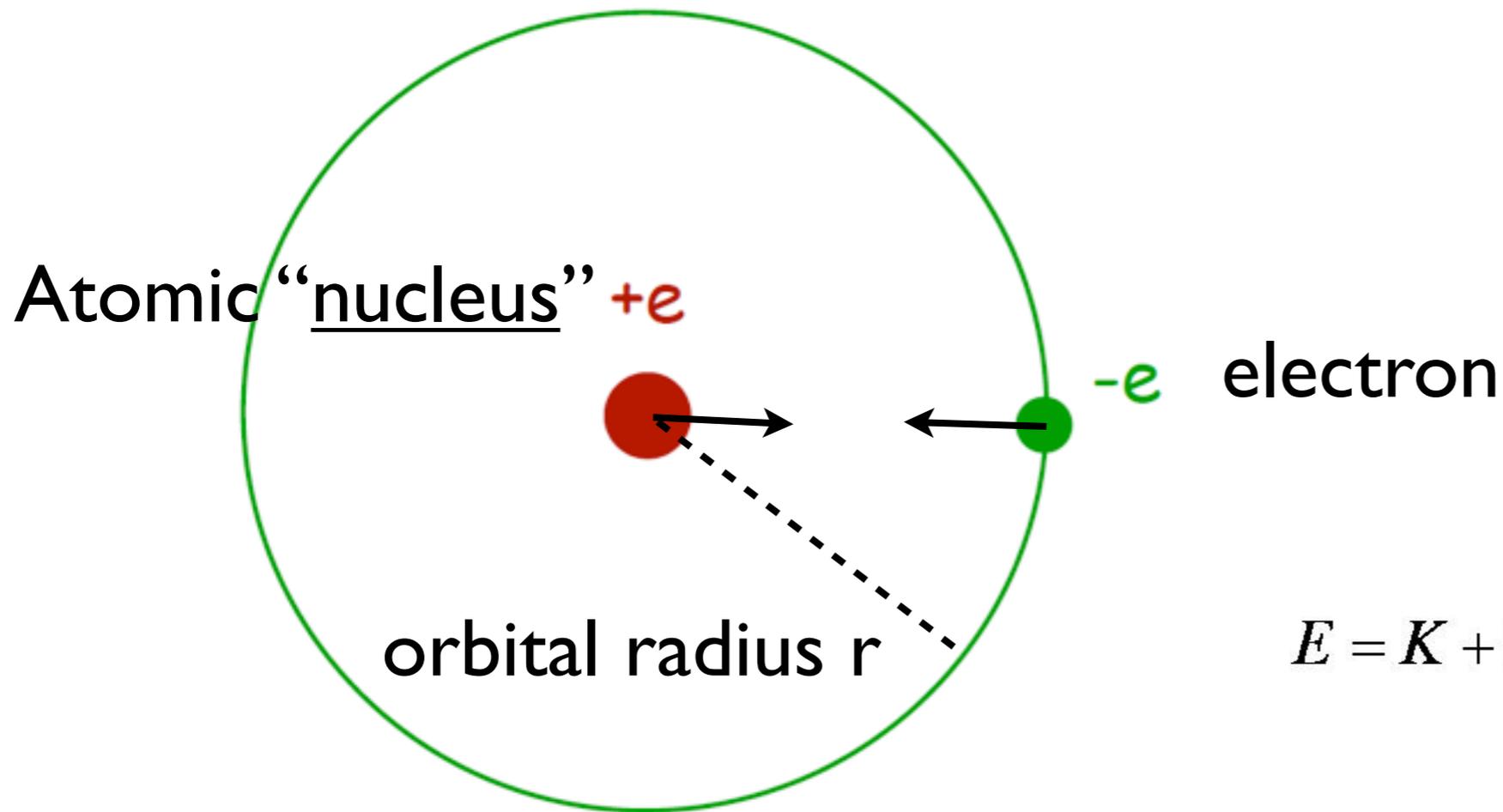
Target:
 atomic density n
 thickness t
 charge $Z_2 e$

Detectors:
 Scattering angle θ
 Distance r

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$$N(\theta) = \frac{N_i n t}{16} \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \frac{Z_1^2 Z_2^2}{r^2 K^2 \sin^4(\theta/2)}$$

Rutherford's Atomic Model



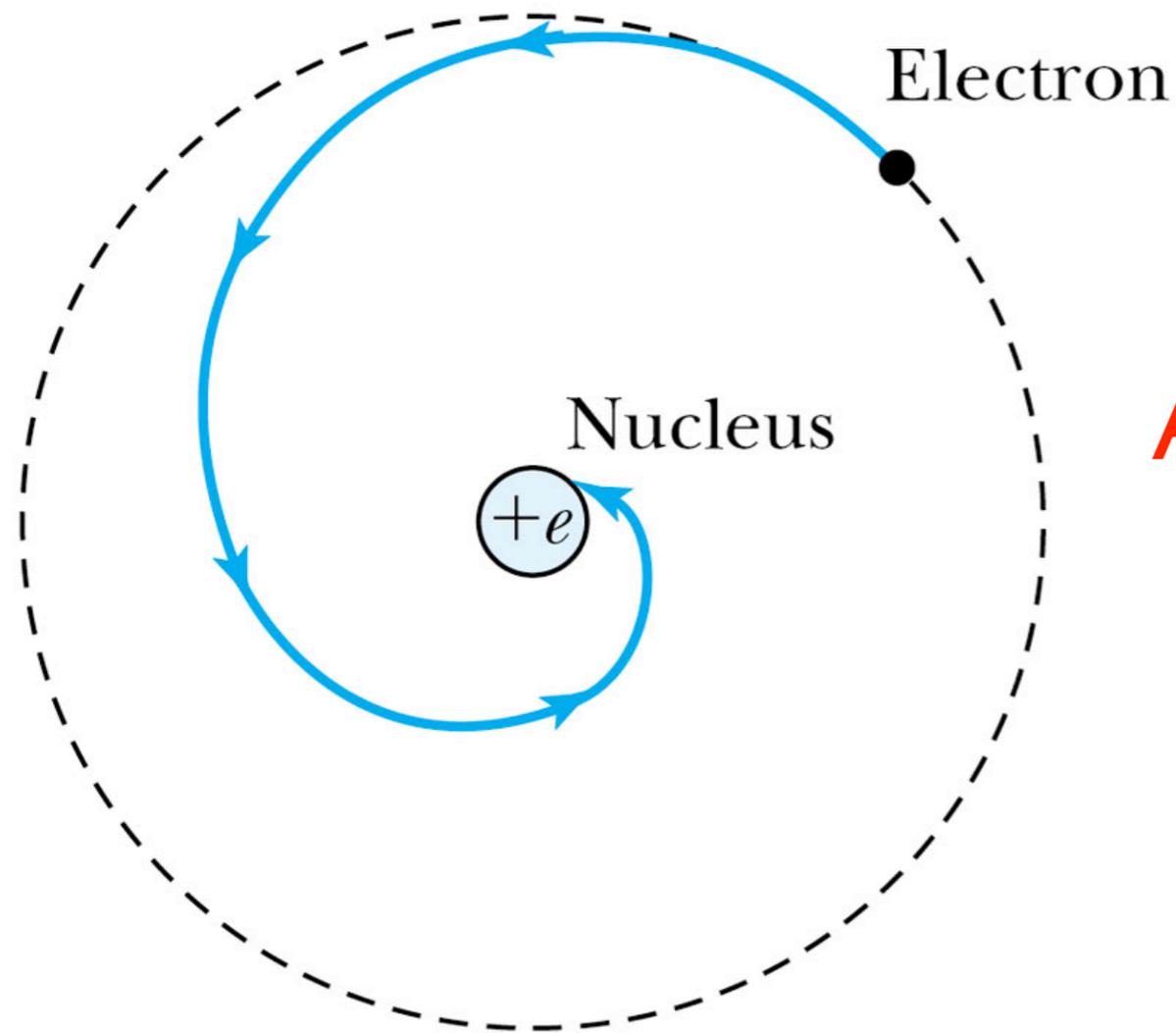
Bound by
Coulomb Force

$$\vec{F}_e = \frac{-1}{4\pi\epsilon_0} \frac{e^2}{r^2} \hat{e}_r = \frac{mv^2}{r}$$

$$E = K + V = \frac{e^2}{8\pi\epsilon_0 r} - \frac{e^2}{4\pi\epsilon_0 r} = \frac{-e^2}{8\pi\epsilon_0 r}$$

nucleus: 10^{-14} m

Problem with the Classical Model



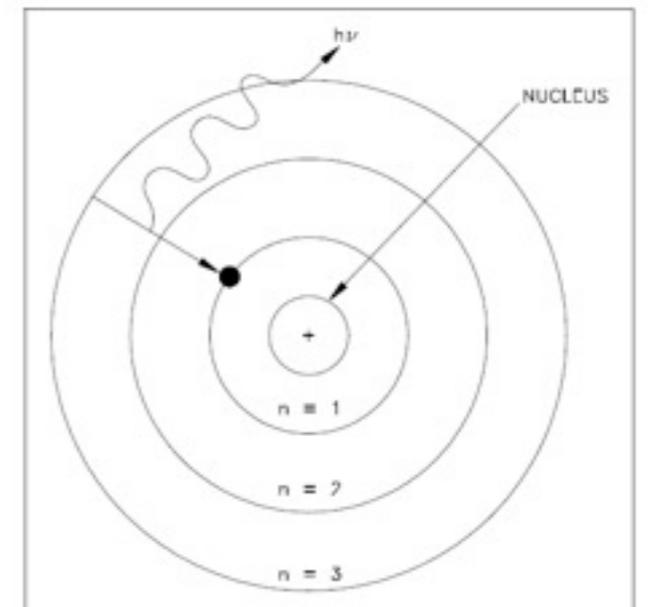
Accelerating electrons should radiate energy - and crash into the nucleus!

Bohr Model

- Electrons can exist only in particular (quantized) “stationary states”.
- Spectral lines correspond to energy *differences* between stationary states, when electrons “jump” between states.
- Angular momentum of electron quantized, equal to an integer multiple of $h/2\pi$!

$$\hbar = h/(2\pi)$$

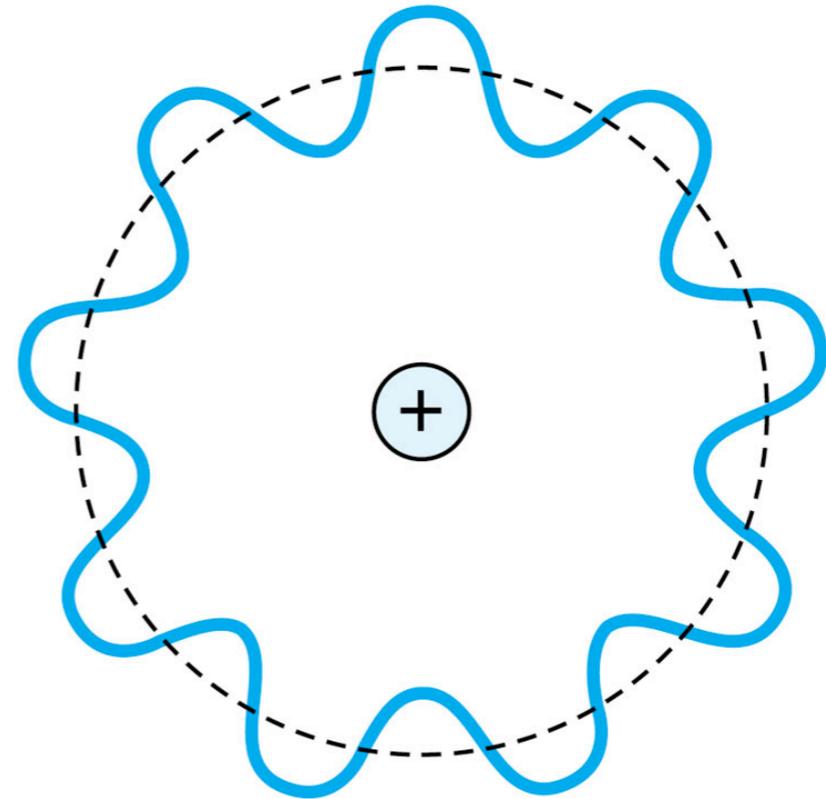
$$\hbar = 1.05 \times 10^{-34} \text{ J}\cdot\text{s}$$



Niels Bohr
1885-1962
Nobel Prize 1922

de Broglie & Bohr

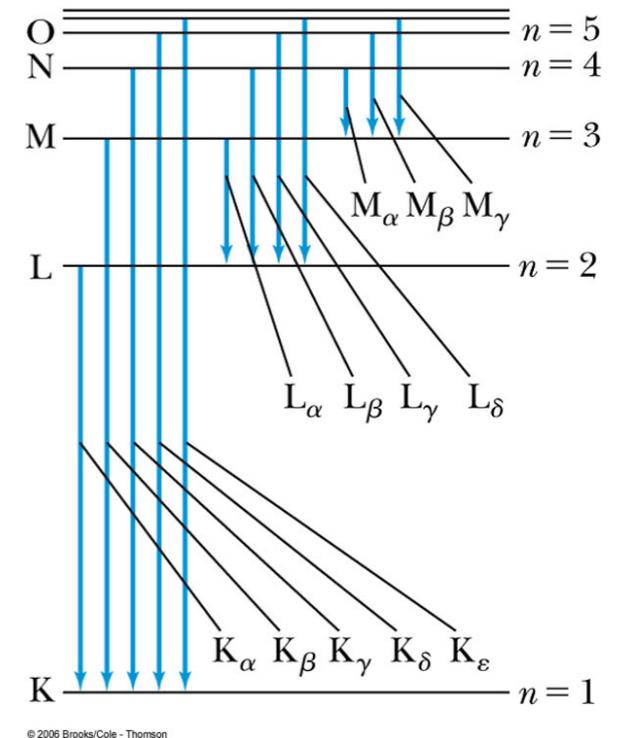
- Bohr: $L = mvr = nh/2\pi$
- de Broglie: $\lambda = h/p$
- $p = mv$, $pr = nh/2\pi$
- Ergo: $n\lambda = 2\pi r$



- “Stationary State” = orbit with an integral number of electron de Broglie wavelengths!
- “Stationary States” = standing electron waves!

Bohr Limitations

- Nucleus is not infinitely heavy
- $m_e \rightarrow \mu = \text{reduced mass}$
- Many electron atoms?
 - Not all electrons in $n=1$ state! Why?
- No systematic way forward
 - What about other systems?
 - How do quantum systems evolve?



Summary, so far

- Bohr's model of the atom
 - Builds on Rutherford's "planetary" model.
 - "Stationery" states = "standing electron waves".
 - Spectral lines = photon emission via electron "jumps" to different levels.

Outline

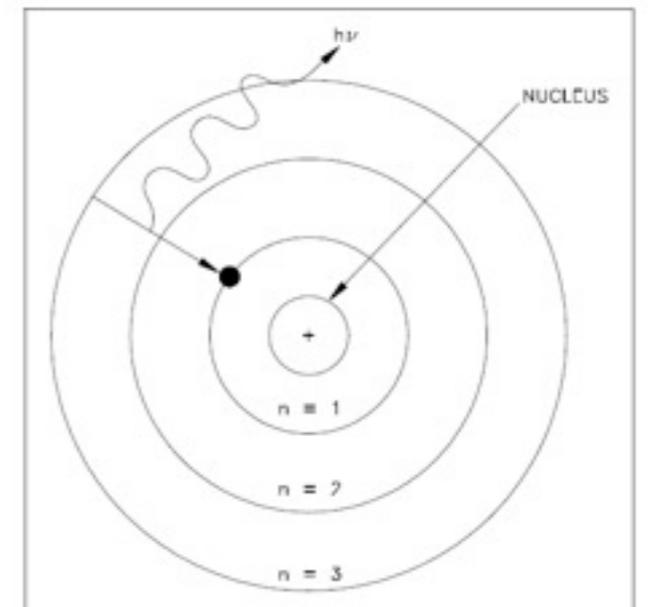
- Bohr Model Numbers
 - Bohr Shell Hypothesis
 - X-ray Spectra and Atomic Number
- Waves vs. Particles
 - Fourier Series/Transforms
 - Complex Exponentials
- Born's Interpretation of Ψ

Bohr Model

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$$\hbar = 1.05 \times 10^{-34} \text{ J}\cdot\text{s}$$



Niels Bohr
1885-1962
Nobel Prize 1922

Bohr Radius (H atom)

Coulomb Force $\frac{e^2}{4\pi\epsilon_0 r_n^2} = \frac{m_e v_n^2}{r_n} = \frac{L^2}{m_e r_n^3}$ ← Centripetal Force

$L = m v_n r_n$

Bohr Quantization $L = n\hbar$

$$r_n = \frac{4\pi\epsilon_0 n^2 \hbar^2}{m_e e^2} = a_0 n^2$$

nucleus: 10^{-14} m (Rutherford)

Bohr Radius $a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = 0.53 \times 10^{-10}$ m = 0.53 \AA

$$\frac{v_n}{c} = \frac{e^2}{4\pi\epsilon_0 c \hbar} \cdot \frac{1}{n} \equiv \frac{\alpha}{n}$$

$$\alpha = \frac{e^2}{4\pi\epsilon_0 c \hbar} \approx \frac{1}{137}$$

Bohr H Energy Levels

$$E_n = -\frac{e^2}{8\pi\epsilon_0 r_n} = -\frac{e^2}{8\pi\epsilon_0 a_0 n^2} \equiv -\frac{E_0}{n^2}$$

$$E_0 \approx 13.6 \text{ eV}$$

n=∞	-----	E = 0
n=4	=====	E = -0.85 eV
n=3	=====	E = -1.51 eV
n=2	-----	E = -3.40 eV

n=1	-----	E = -13.6 eV
-----	-------	--------------

Transition from state n to state m:

$$\begin{aligned}
 E_{k \rightarrow n} &= E_k - E_n & \frac{1}{\lambda} &= \frac{E_0}{hc} \left(\frac{1}{n^2} - \frac{1}{k^2} \right) \\
 &= \frac{hc}{\lambda} & &= R_\infty \left(\frac{1}{n^2} - \frac{1}{k^2} \right) \\
 &= E_0 \left(\frac{1}{n^2} - \frac{1}{k^2} \right) & R_\infty &= \frac{13.6 \text{ eV} \cdot 1.602 \times 10^{-19} \text{ J/eV}}{1.986 \times 10^{-25} \text{ J m}} \\
 & & &= 1.097 \times 10^7 \text{ m}^{-1}
 \end{aligned}$$

Bohr model reproduces Rydberg formula!

Reduced Mass Correction

(Replacing m_e by μ_e)

The electron and hydrogen nucleus actually revolve about their mutual center of mass as shown in Figure 4.17. This is a two-body problem, and our previous analysis should be in terms of r_e and r_M instead of just r . A straightforward analysis derived from classical mechanics shows that this two-body problem can be reduced to an equivalent one-body problem in which the motion of a particle of mass μ_e moves in a central force field around the center of mass. The only change required in the results of Section 4.4 is to replace the electron mass m_e by its **reduced mass** μ_e where

$$\frac{1}{\mu_e} = \frac{1}{m_e} + \frac{1}{M} \Rightarrow \mu_e = \frac{m_e M}{m_e + M} = \frac{m_e}{1 + \frac{m_e}{M}} \quad (4.36)$$

and M is the mass of the nucleus (see Problem 53). In the case of the hydrogen atom, M is the proton mass, and the correction for the hydrogen atom is $\mu_e = 0.999456 m_e$. This difference can be measured experimentally. The Rydberg constant for infinite nuclear mass, R_∞ , defined in Equation (4.29), should be replaced by R , where

$$R = \frac{\mu_e}{m_e} R_\infty = \frac{1}{1 + \frac{m_e}{M}} R_\infty = \frac{\mu_e e^4}{4\pi c \hbar^3 (4\pi\epsilon_0)^2}$$

(QED)

$\frac{\Delta R}{R} \sim 10^{-12}$

(4.37)

We can use the precision measurement of R to determine M via the ratio $\frac{m_e}{M}$.

The Rydberg constant for hydrogen is $R_H = 1.096776 \times 10^7 \text{ m}^{-1}$.

Concept Test

- Let's apply the Bohr model to carbon, whose nucleus has $Z=+6$. The number of electrons in the neutral atom is therefore
 - 4
 - 5
 - 6  Are all six electrons in the lowest energy stationary state?
 - 7

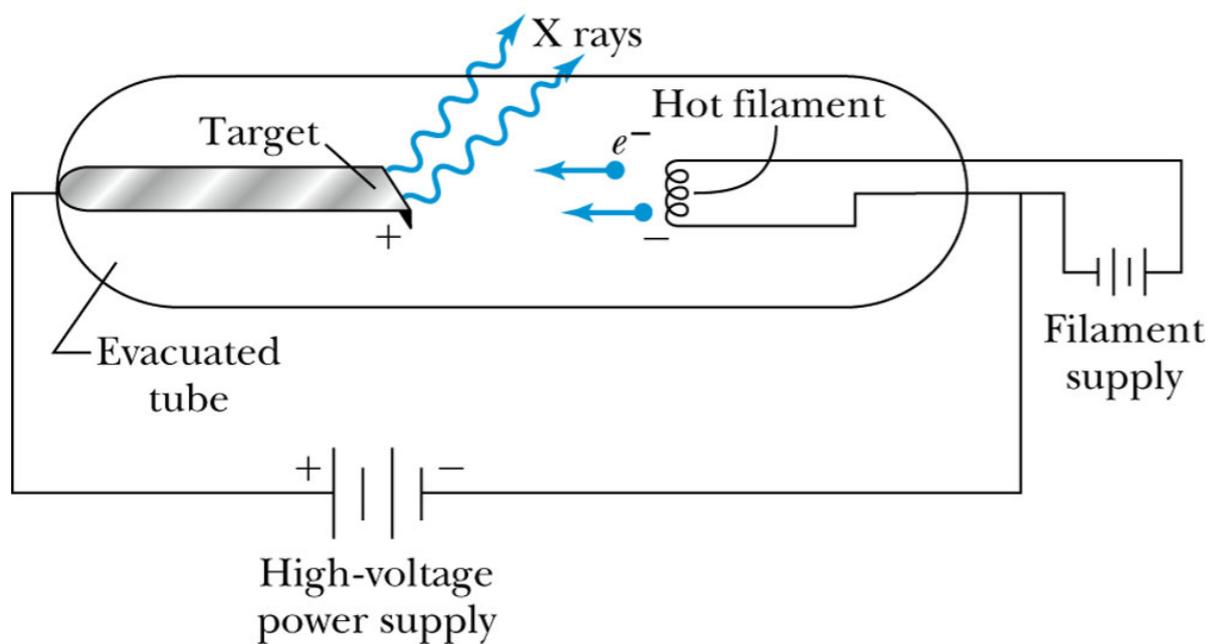
Bohr Shell Hypothesis

$$\frac{1}{\lambda_{k \rightarrow n}} = Z^2 R \left(\frac{1}{n^2} - \frac{1}{k^2} \right)$$

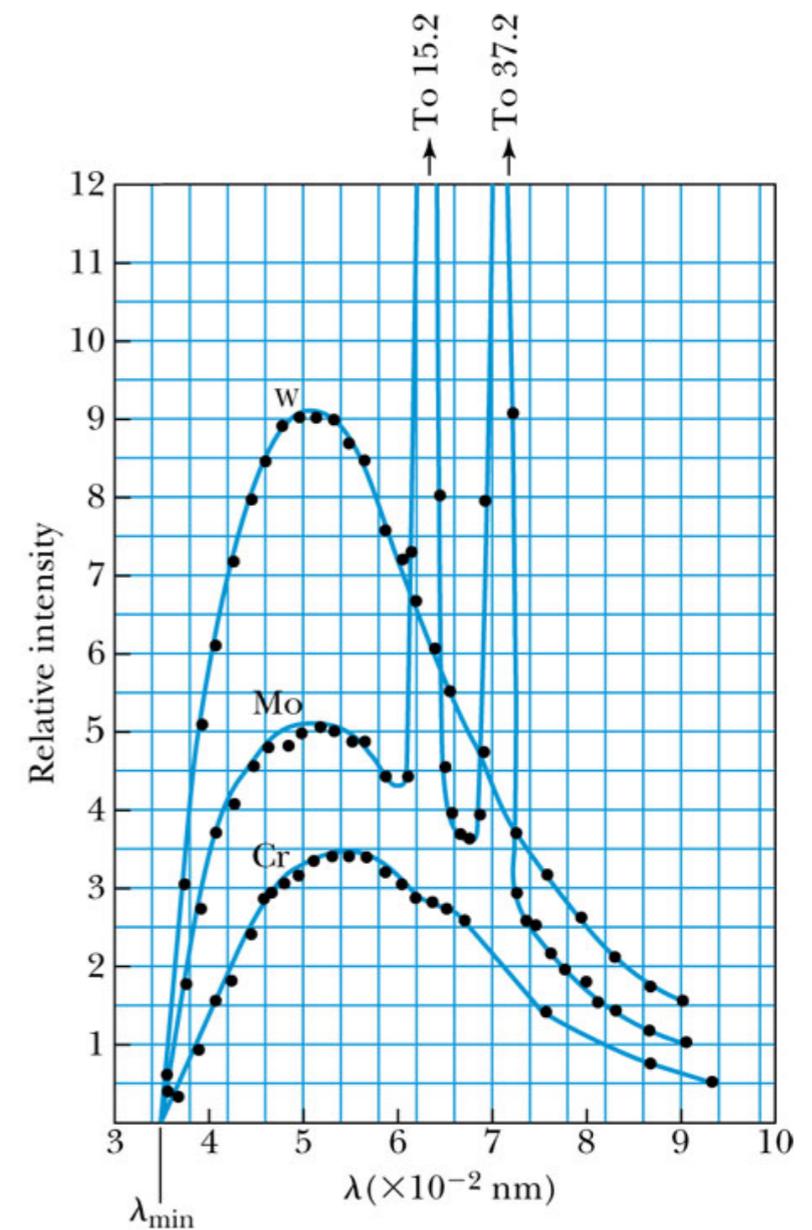
- Bohr model generalizes to any single-electron atom: $e^2 \rightarrow Ze^2$
- Bohr model yields many quantized energy levels for any atom, depending on n .
- Bohr asserted that any given shell could only hold a certain number of electrons - after it was filled, electrons must occupy the next available level. **Why? (We will see!)**

X-Ray Spectra Peaks:

Inverse photoelectric effect:

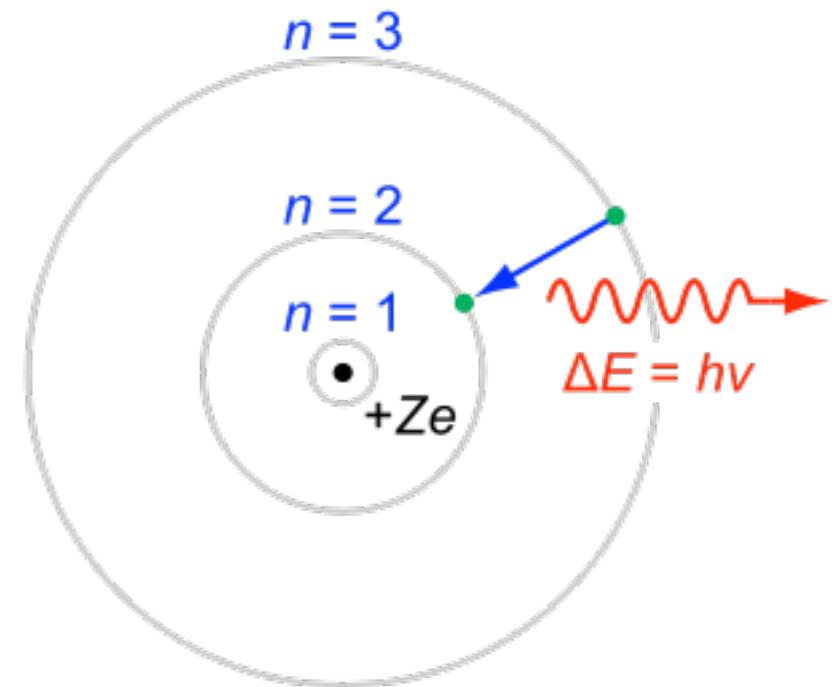


Can be explained in terms of shell hypothesis...

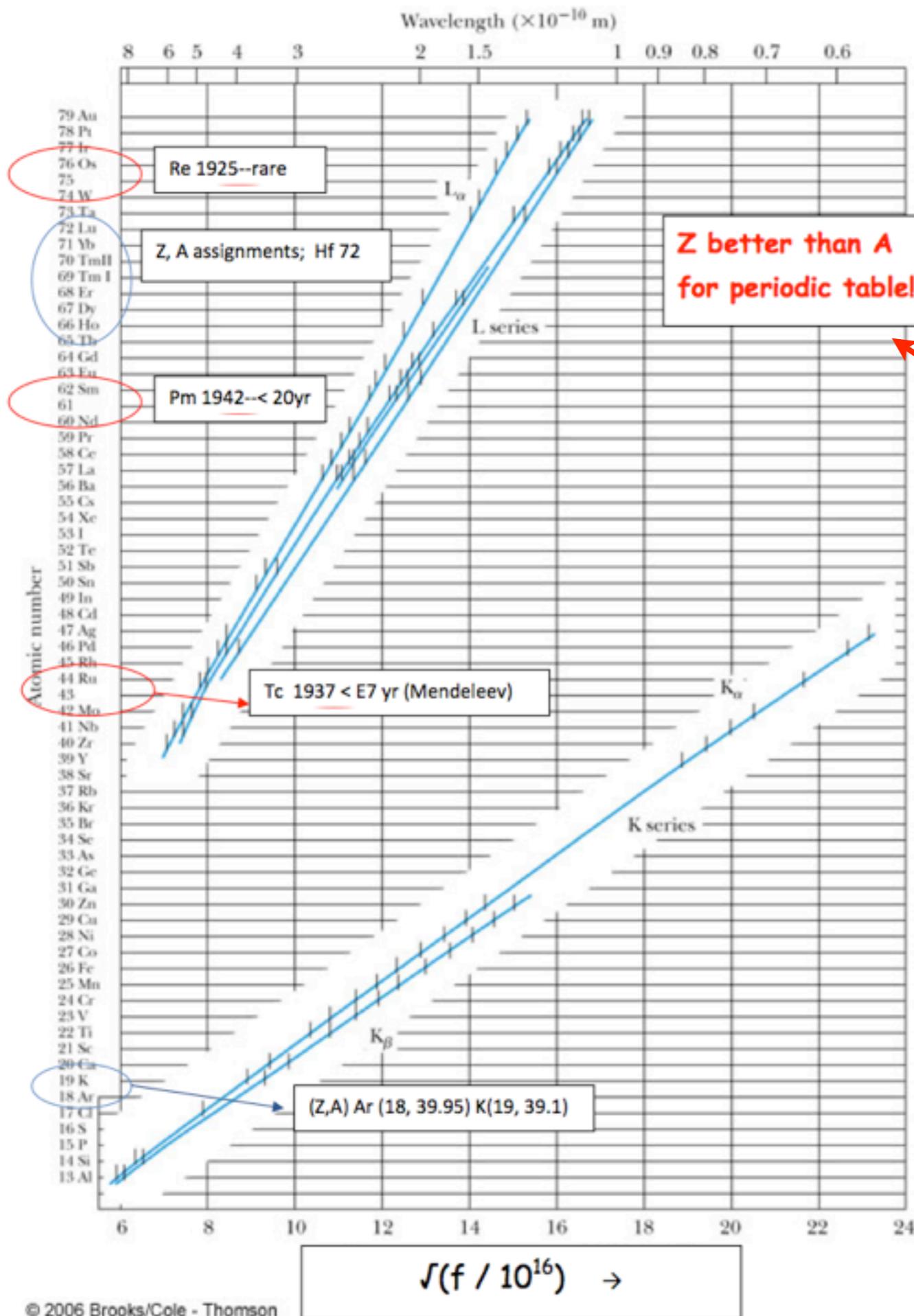


X-ray Peaks II

- X-rays excite electrons from $n=1$ “shell” (K-shell).
- An electron from an upper shell cascades down to take its place
- Since $E_n \propto Z^2$, we expect square root of peak frequencies to be linear in Z !



Moseley Plot (1913)



Z better than A for periodic table!

Corrected Mendeleev!



Henry Moseley
1887-1915