

The Atomic states

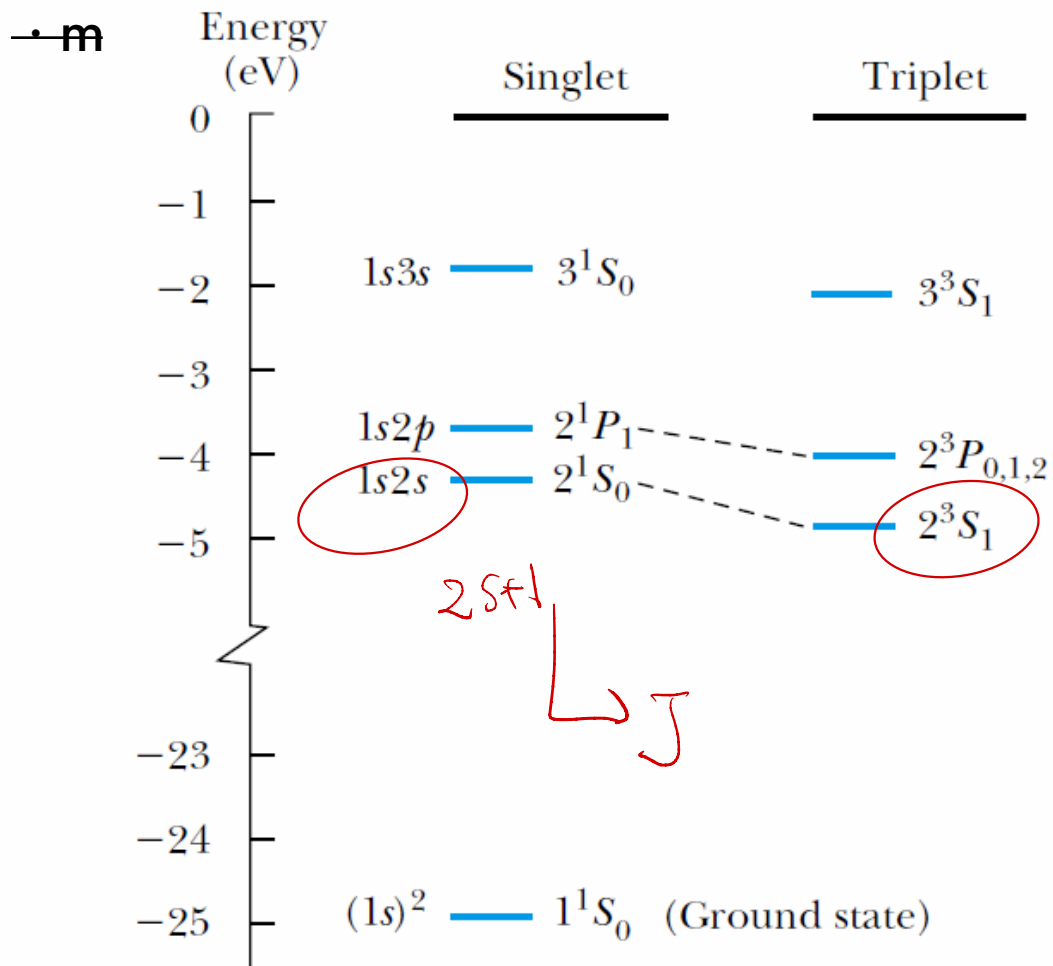


Figure 8.13 The low-lying atomic states of helium are shown. The ground state (1^1S_0) is some 20 eV below the grouping of the lowest excited states. The level indicated by $3^3P_{0,1,2}$ is actually three states (3^3P_0 , 3^3P_1 , 3^3P_2), but the separations are too small to be indicated.

2^3S_1 is allowed because $1s$ and $2s$ are in different n shells, so that they are distinguishable states.

Two electrons $1s2s$

① $l_1=0, s_1=\frac{1}{2}$
 ② $l_2=0, s_2=\frac{1}{2}$

$\vec{L} = \vec{L}_1 + \vec{L}_2$; $\vec{S} = \vec{S}_1 + \vec{S}_2$
 ($l=0$) ; ($s=0, 1$)

$\vec{J} = \vec{L} + \vec{S}$
 ($J=0, 1$)

2^1S_0 & 2^3S_1

The Atomic states

Spectroscopic Notation



Hydrogen energy levels

Less than thermal energy

$$4.5 \times 10^{-5} \text{ eV} = \frac{3}{2} kT \text{ at } 0.35 \text{ K}$$

Barely resolvable since the Doppler broadening is about 4 GHz

$$\Delta E = 4.5 \times 10^{-5} \text{ eV}$$

$$\nu = 10.9 \text{ GHz}$$

$$\lambda = 2.8 \text{ cm}$$

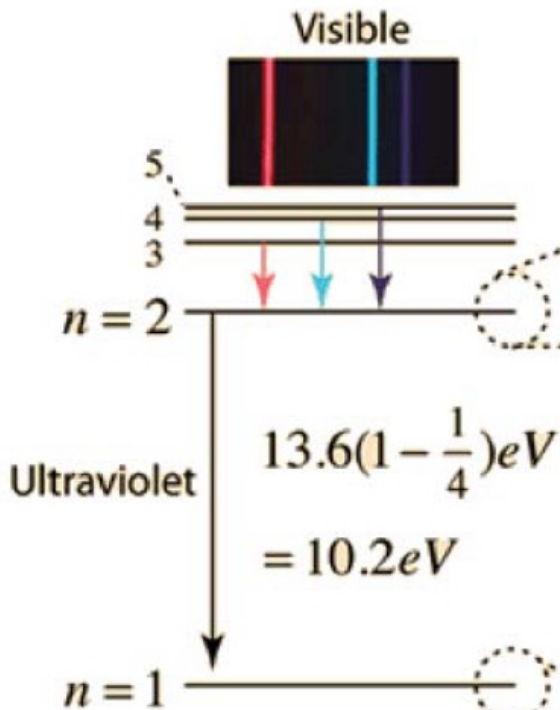
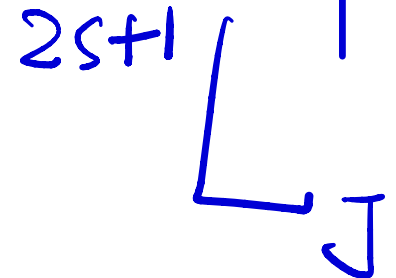
$2p_{3/2}$ (for $s = \frac{1}{2}$, $l = 1$)
 $2p_{1/2}$

$$\Delta E = 4.372 \times 10^{-6} \text{ eV}$$

$$\nu = 1.0576 \text{ GHz}$$

$$\lambda = 28.37 \text{ cm}$$

Lamb shift



$$E_n = \frac{-13.6z^2}{n^2} \text{ eV}$$

$$\Delta E = 5.9 \times 10^{-6} \text{ eV}$$

$$\nu = 1.42 \text{ GHz}$$

$$\lambda = 21 \text{ cm}$$

Hyperfine structure

N = nuclear spin
 E = electron spin

x 100,000

x 100,000

Fine structure

$2s_{1/2}$
 $2p_{1/2}$

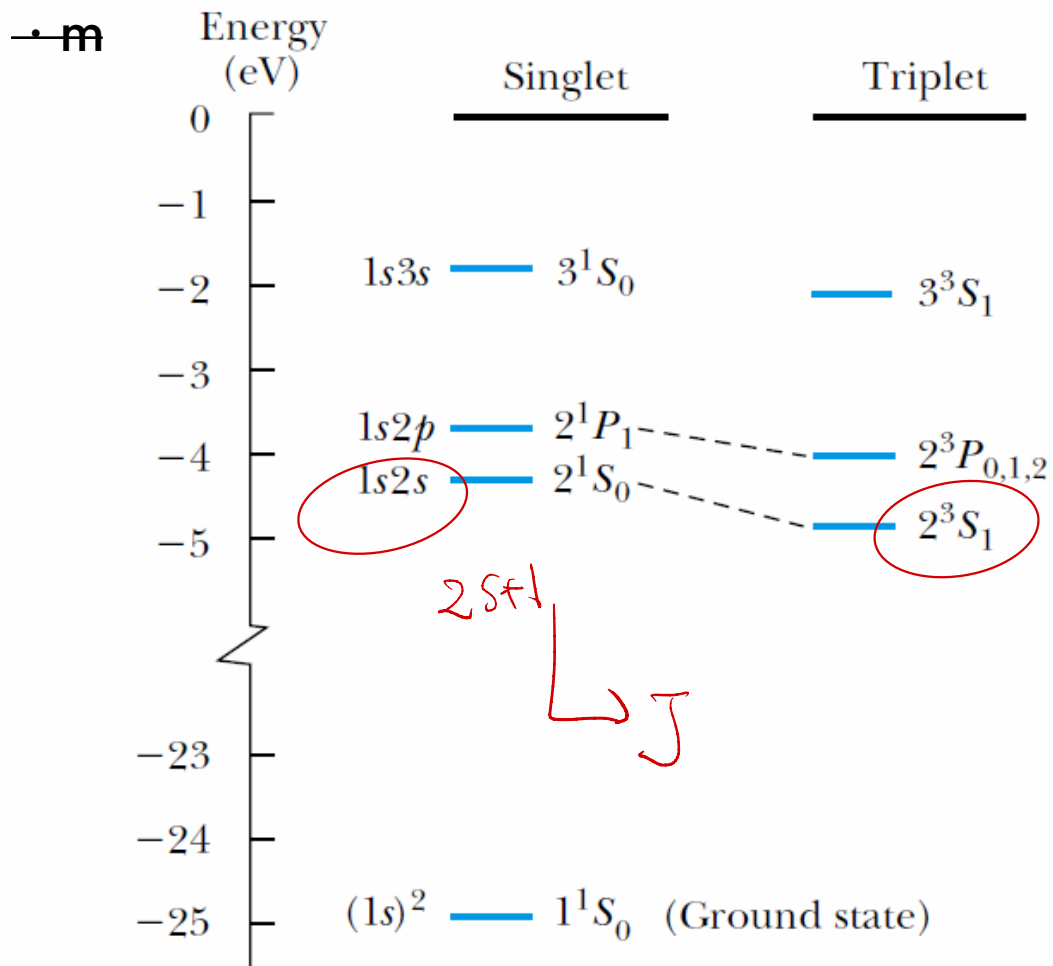


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Two electrons 1s2s

① $l_1=0, s_1=\frac{1}{2}$
 ② $l_2=0, s_2=\frac{1}{2}$
 $\Rightarrow \vec{L} = \vec{L}_1 + \vec{L}_2$; $\vec{S} = \vec{S}_1 + \vec{S}_2$
 ($l=0$) ; ($s=0, 1$)
 $\vec{J} = \vec{L} + \vec{S}$
 ($J=0, 1$)
 $\Rightarrow 2^1S_0 \text{ and } 2^3S_1$

What are the possible energy states for atomic carbon?

Strategy The element carbon has two $2p$ subshell electrons outside the closed $2s^2$ subshell. Both electrons have $\ell = 1$, so we have $L = 0, 1$, or 2 using the LS coupling scheme. The spin angular momentum is $S = 0$ or 1 .

$S=0$ (antisymmetric) $\frac{1}{\sqrt{2}}(\uparrow\downarrow - \downarrow\uparrow)$
 $S=1$ (symmetric) spin wave function $\uparrow\downarrow, \frac{1}{\sqrt{2}}(\uparrow\downarrow + \downarrow\uparrow), \downarrow\downarrow$

$2S+1$
 L, J
 $\vec{L} = \vec{L}_1 + \vec{L}_2$

Spatial (orbital) wave function
 $(-1)^L = \begin{cases} -1 & \text{antisymmetric} \\ 1 & \text{symmetric} \end{cases}$

S	L	J	Spectroscopic Notation
0	0	0	1S_0
	1	1	1P_1 $S=0, L=1$ Not allowed
	2	2	1D_2
1	0	1	3S_1 $S=1, L=0$ Not allowed
	1	0, 1, 2	$^3P_{0,1,2}$
	2	1, 2, 3	$^3D_{1,2,3}$ $S=1, L=2$ Not allowed

For fermions (such as electrons),
 Total wave function
 $\equiv (\text{Spin wave func.}) * (\text{spatial, orbital wave func.})$
 $= \text{anti-symmetric}$
 (Pauli Exclusion Principle)

\Rightarrow Require

$(-1)^{S+1} \cdot (-1)^L = -1$

(for the LS coupling scheme) $\frac{1}{\sqrt{2}}(\uparrow\downarrow - \downarrow\uparrow)$

The $4p$ and $4d$ orbitals are distinguishable quantum states, for having different angular momentum quantum numbers.

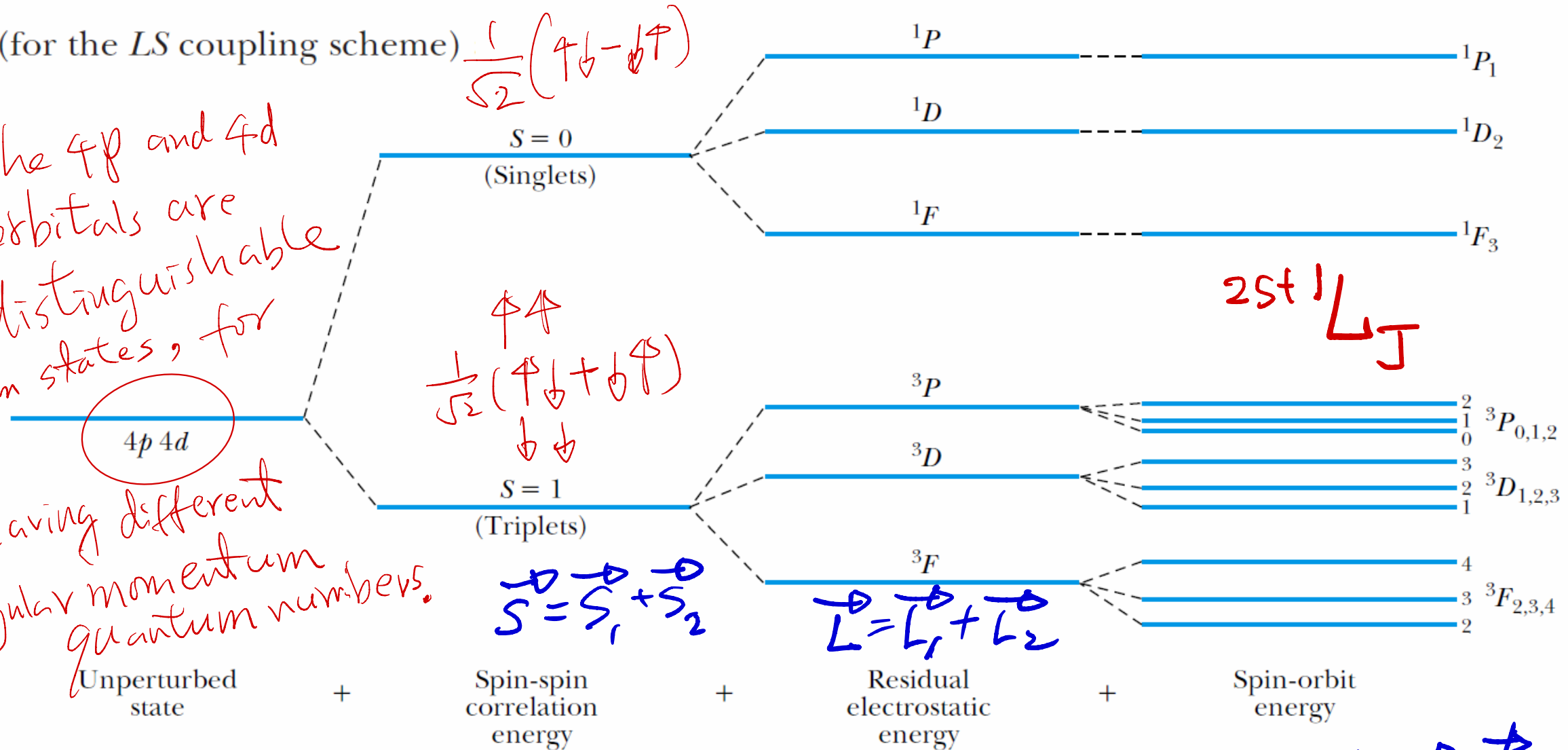


Figure 8.9 Schematic diagram indicating the increasing fine-structure splitting due to different effects. This case is for an atom having two valence electrons, one in the $4p$ and the other in the $4d$ state. The energy is not to scale. From R. B. Leighton, *Principles of Modern Physics*, New York: McGraw-Hill (1959), p. 261. Used with permission.

Transition Matrix element
 $\langle f | H_{int} | i \rangle$
 $= \langle f | (-e \vec{r} \cdot \vec{E}) | i \rangle$
 for the electric dipole interaction

	Singlets				Triplets			
Electron configuration \rightarrow	$3sns$	$3snp$	$3snd$	$3snf$	$3sns$	$3snp$	$3snd$	$3snf$
Term symbol \rightarrow	1S_0	1P_1	1D_2	1F_3	3S_1	$^3P_{0,1,2}$	$^3D_{1,2,3}$	$^3F_{2,3,4}$

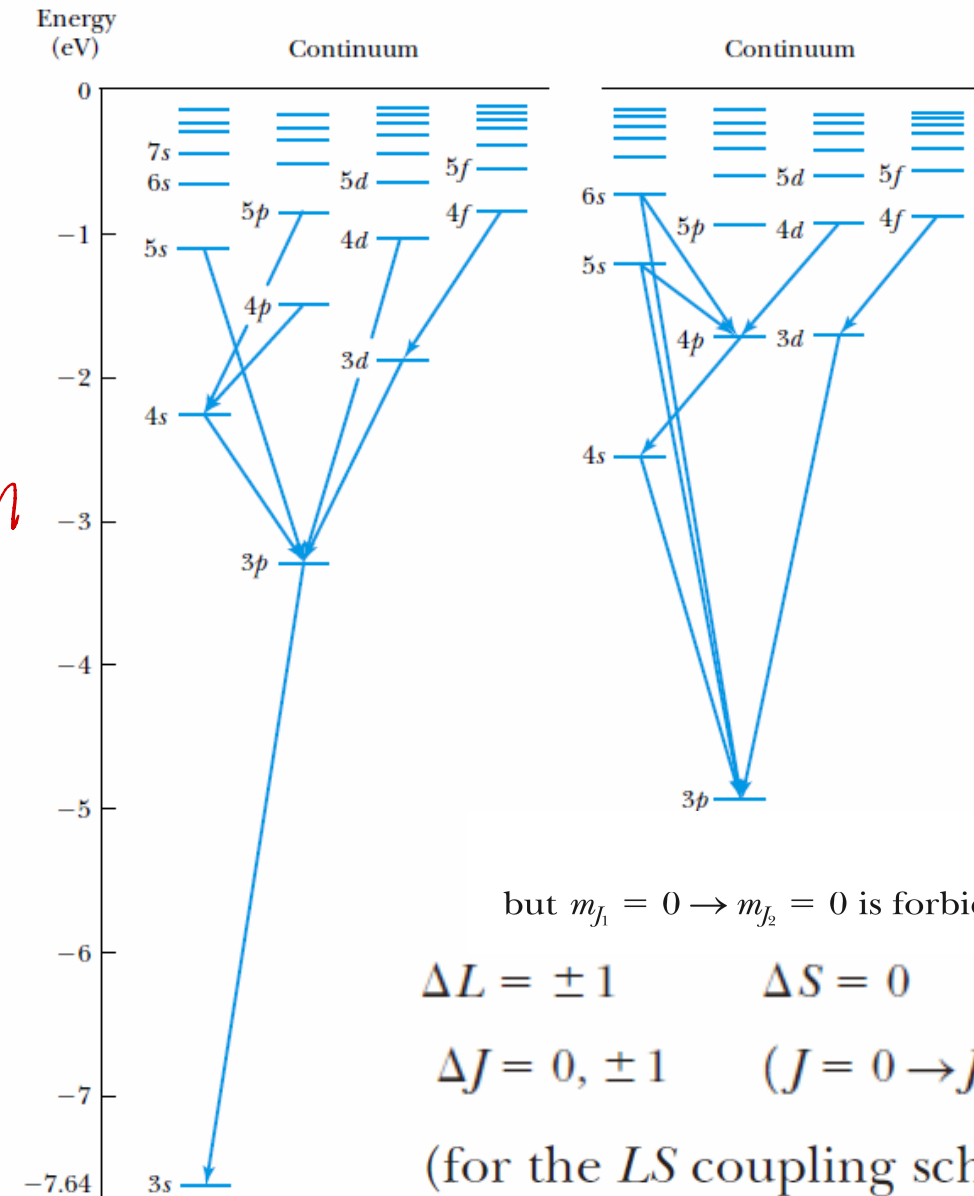


Figure 8.10 Energy-level diagram for magnesium (two-electron atom) with one electron in the $3s$ subshell and the other electron excited into the nl subshell indicated. The singlet and triplet states are separated, because transitions between them are not allowed by the $\Delta S = 0$ selection rule. Several allowed transitions are indicated.

$\Delta m_j = \pm 1, 0$
 but $m_{j_1} = 0 \rightarrow m_{j_2} = 0$ is forbidden when $\Delta J = 0$.

$\Delta L = \pm 1$ $\Delta S = 0$
 $\Delta J = 0, \pm 1$ ($J = 0 \rightarrow J = 0$ is forbidden)

(for the LS coupling scheme)

What are the possible energy states for atomic carbon?

Strategy The element carbon has two $2p$ subshell electrons outside the closed $2s^2$ subshell. Both electrons have $\ell = 1$, so we have $L = 0, 1, \text{ or } 2$ using the LS coupling scheme. The spin angular momentum is $S = 0$ or 1 .

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$2S+1$
 L
 J
 $\vec{L} = \vec{L}_1 + \vec{L}_2$

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⇒ Require

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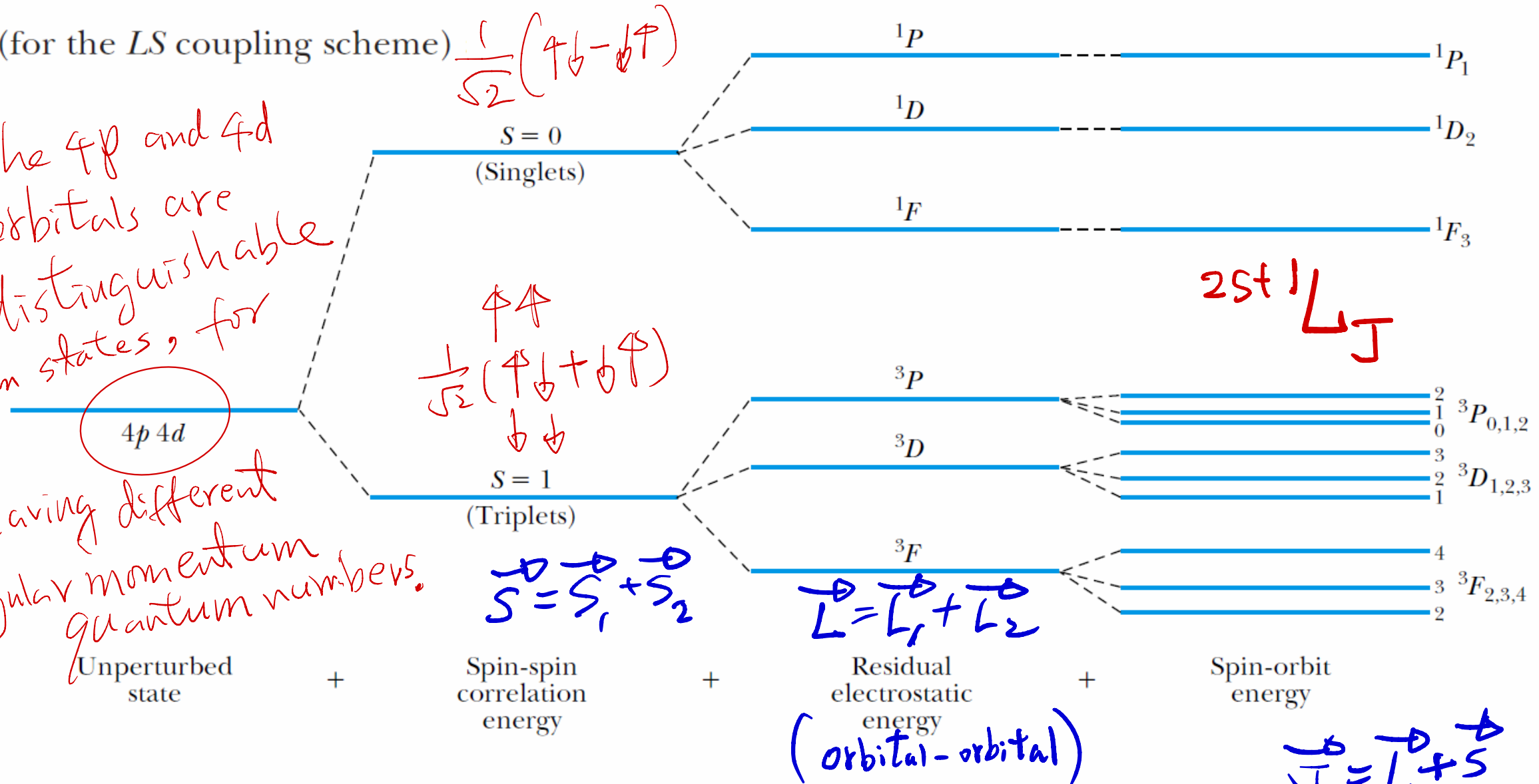
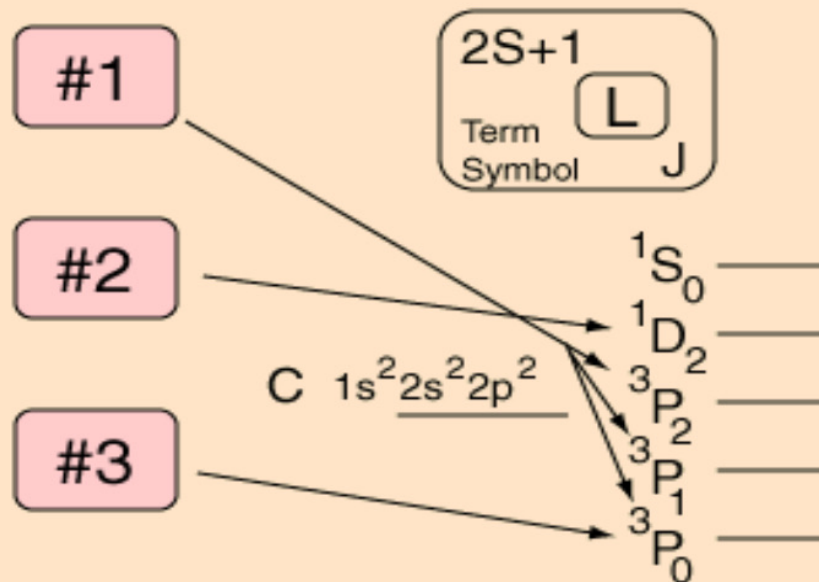


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Hund's Rules

1. The term with maximum multiplicity lies lowest in energy.
2. For a given multiplicity, the term with the largest value of L lies lowest in energy.
3. For atoms with less than half-filled shells, the level with the lowest value of J lies lowest in energy.

Click on the indicated buttons for more detail.



Hund's rules assume combination to form S and L, or imply [L-S \(Russell-Saunders\) coupling](#).

Note: Some references, such as [Haken & Wolf](#), use Hund's Rule #1 to apply to the nature of full shells and subshells. Full shells and subshells contribute nothing to the total angular momenta L and S. If you call this Hund'e Rule #1, then the above rules will be bumped up one in number. I don't know which is the more common practice.

[Exceptions to Hund's Rules](#)

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[Atomic structure concepts](#)

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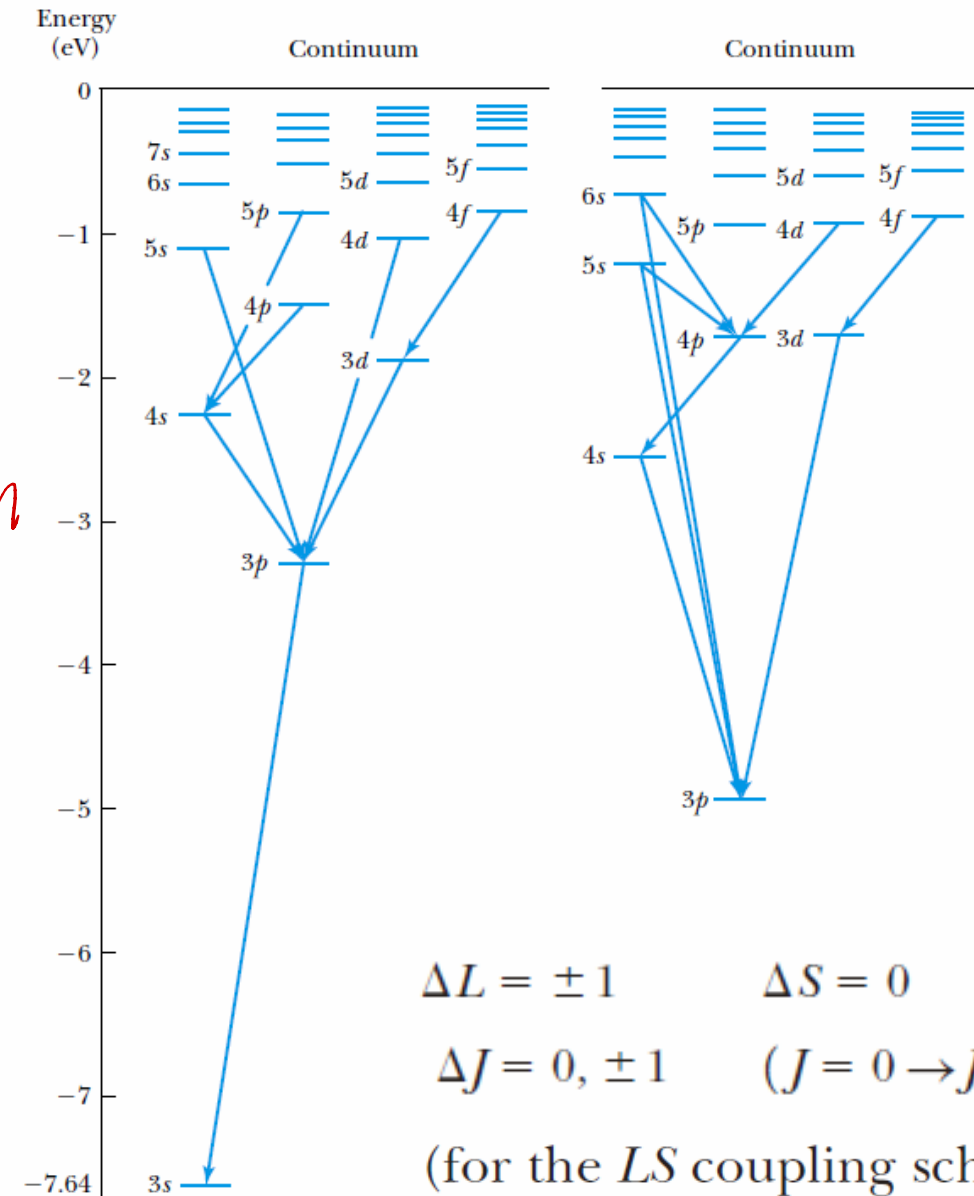
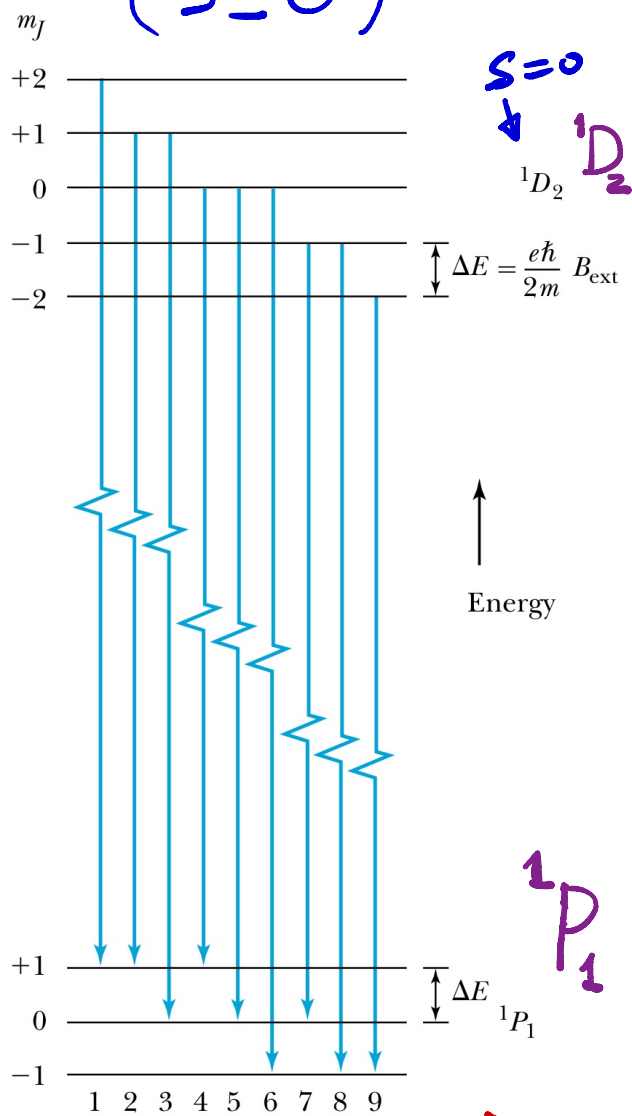


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$\Delta L = \pm 1$ $\Delta S = 0$
 $\Delta J = 0, \pm 1$ ($J = 0 \rightarrow J = 0$ is forbidden)
 (for the LS coupling scheme)

normal Zeeman Effect ($s=0$)



with \vec{B}_{ext}

anomalous Zeeman Effect ($s \neq 0$)

